

A cancellation mechanism for dark matter-nucleon interaction: non-Abelian case

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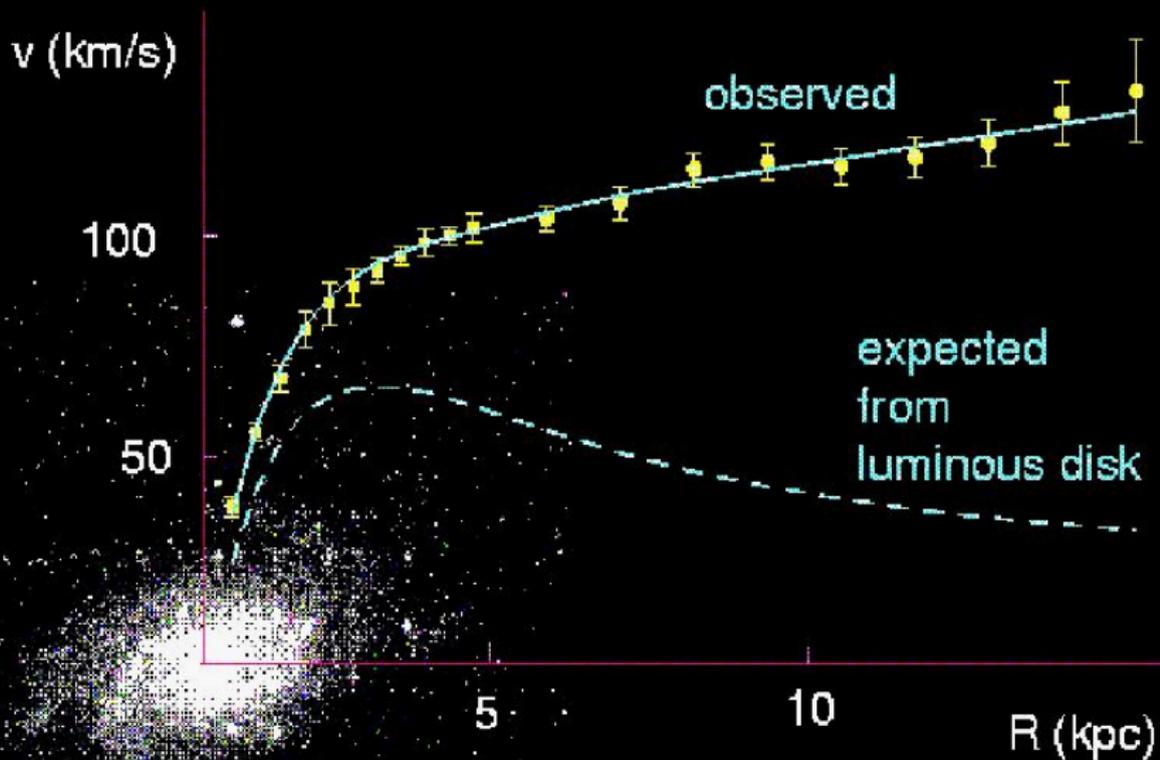
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In collaboration with:

**Christian Gross, Alexandros Karam, Oleg Lebedev, Kyriakos
Tamvakis**

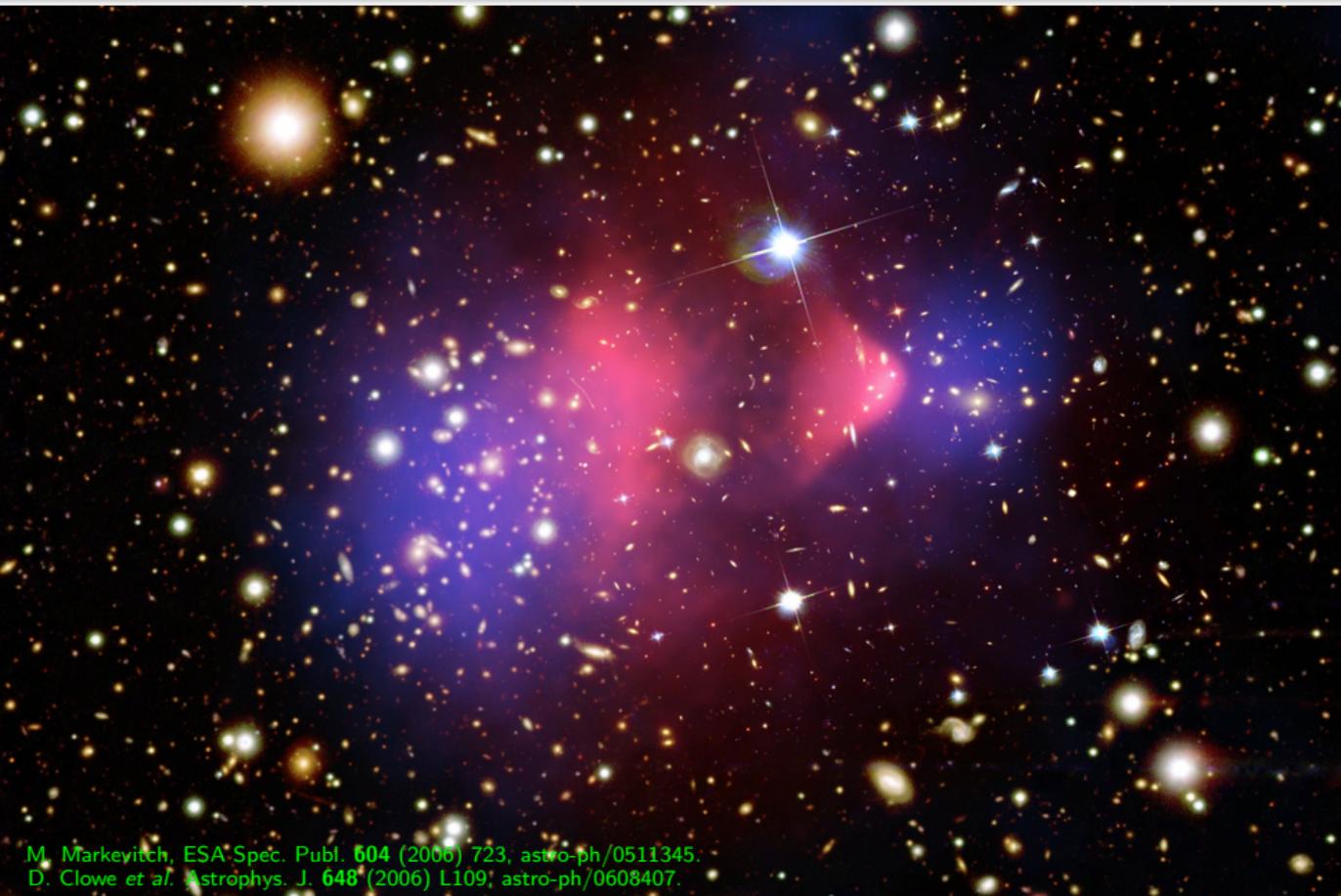
Introduction

Velocity distribution



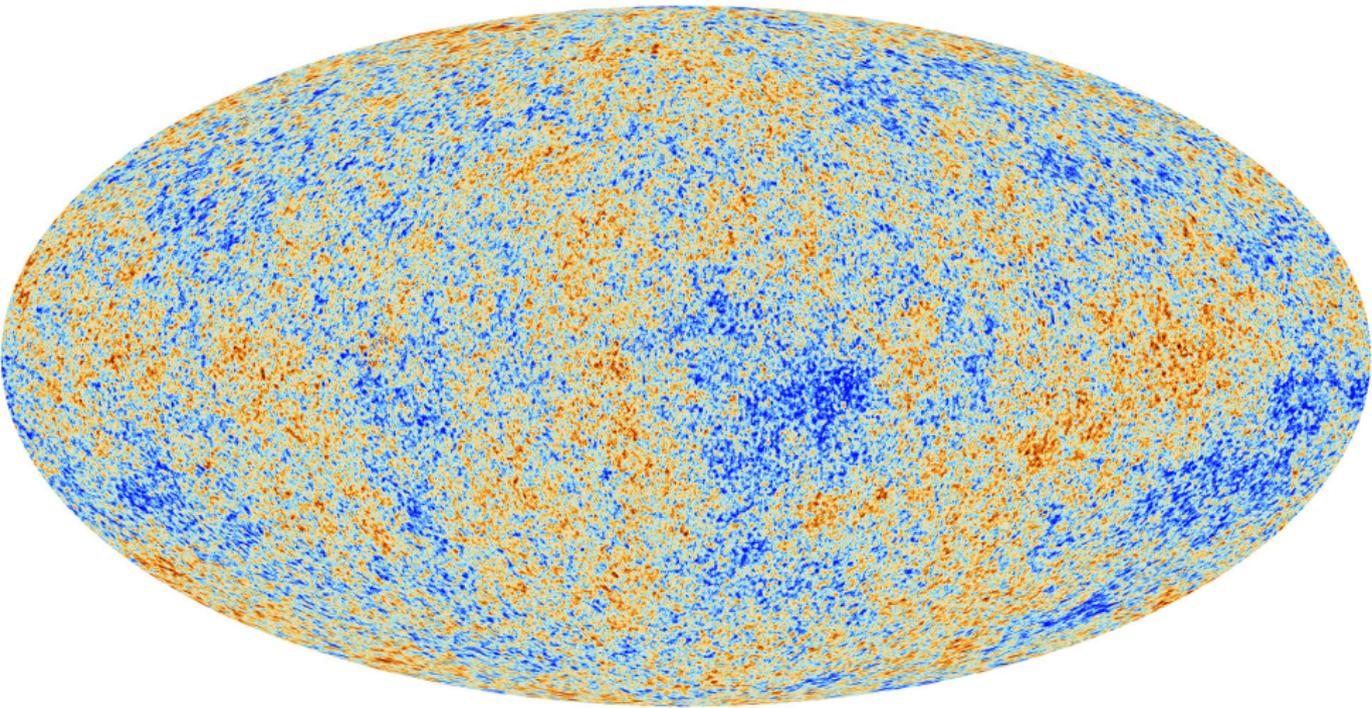
E. Corbelli and P. Salucci, *Mon. Not. Roy. Astron. Soc.* **311** 441 (2000), arXiv:astro-ph/9909252.

Bullet Cluster



M. Markevitch, *ESA Spec. Publ.* **604** (2006) 723; [astro-ph/0511345](#).
D. Clowe *et al.* *Astrôphys. J.* **648** (2006) L109; [astro-ph/0608407](#).

Cosmic Microwave Background



Dark Matter relic abundance¹ $\Omega_{DM}h^2 \approx 0.12$.

¹P. A. R. Ade et al. [Planck Collaboration] *Astron. Astrophys.* **594** (2016) A13, arXiv:1502.01589.

Some general characteristics

- Stable or very slow decay rate.
- Electrically Neutral.
- Cold/Warm and non-relativistic today.
- Smaller allowed mass² $m_{DM} \gtrsim \mathcal{O}(1 - 10 \text{ keV})$.
- Non-Baryonic.

²V. Iršič et al. Phys. Rev. D **96** (2017) no.2, 023522, arXiv:1702.01764.

Three stages:

- Equilibrium (production \leftrightarrow annihilation), $T \gtrsim m_{DM}$.
- Production stops, $T \lesssim m_{DM}$.
- Annihilation stops, $T = T_{FO} \approx \frac{m_{DM}}{25}$.

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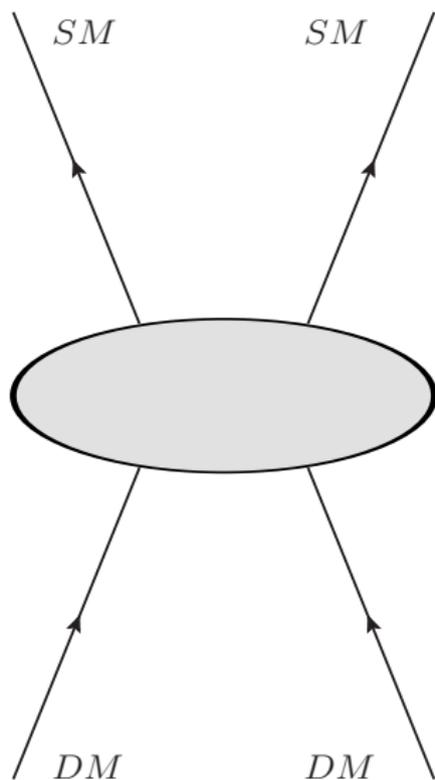
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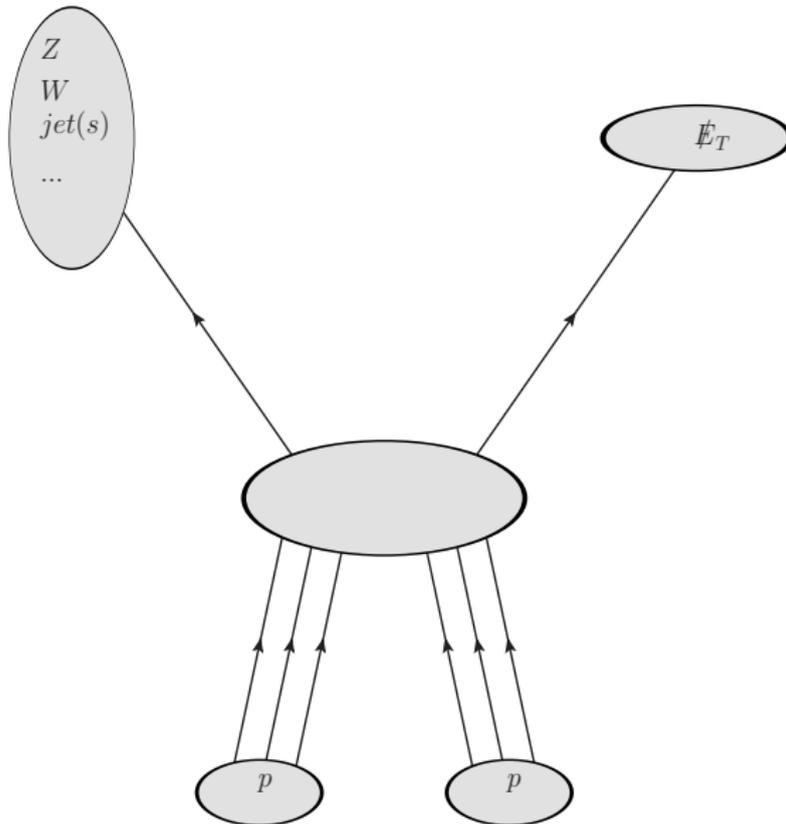
WIMP= Weakly Interacting Massive Particle.

$$m_{DM} \sim 100 \text{ GeV} \Rightarrow \Omega_{MD} h^2 \sim 0.1 \frac{10^{-8} \text{ GeV}^{-2}}{\langle \sigma v \rangle} .$$

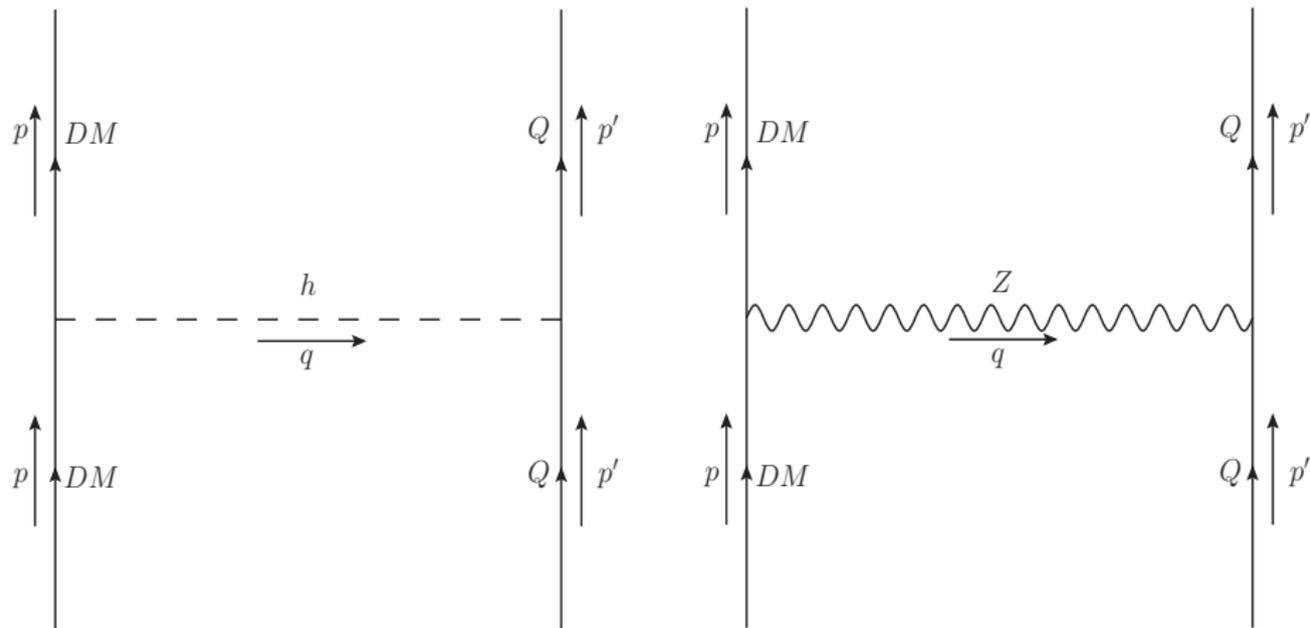
Typical mass scale of weak interactions \leftrightarrow Typical cross section of weak interactions (*WIMP-miracle*).

Indirect detection





Direct detection



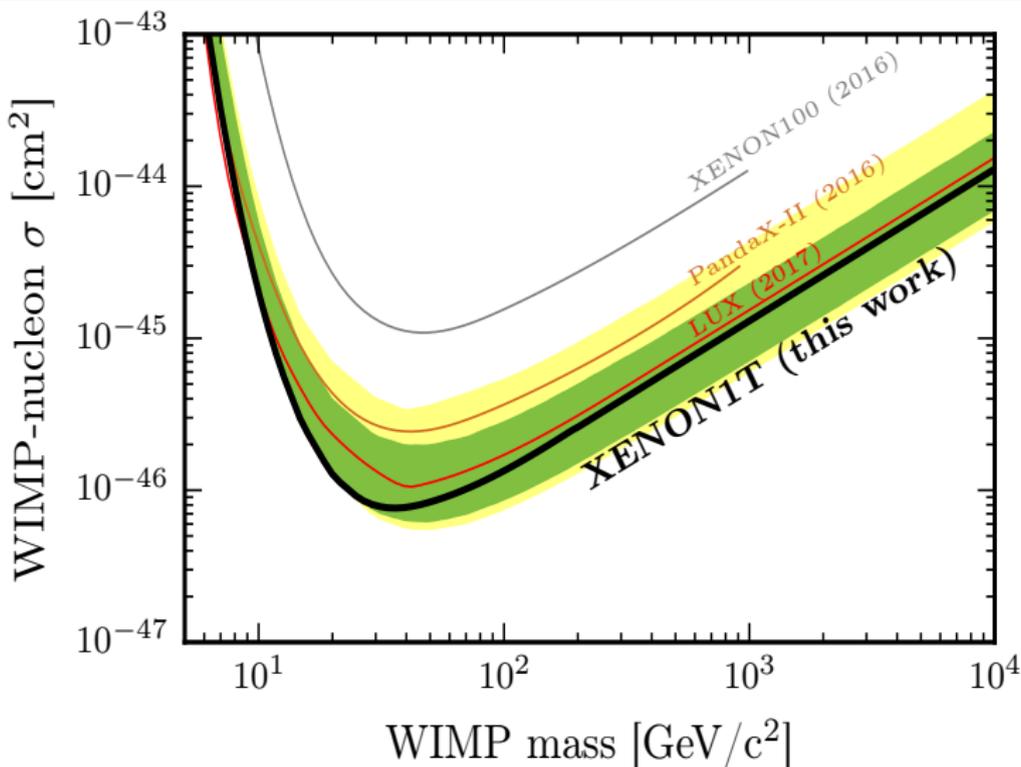


Figure : E. Aprile *et al.* [XENON Collaboration], Phys. Rev. Lett. **119**, no. 18, 181301 (2017), arXiv:1705.06655.

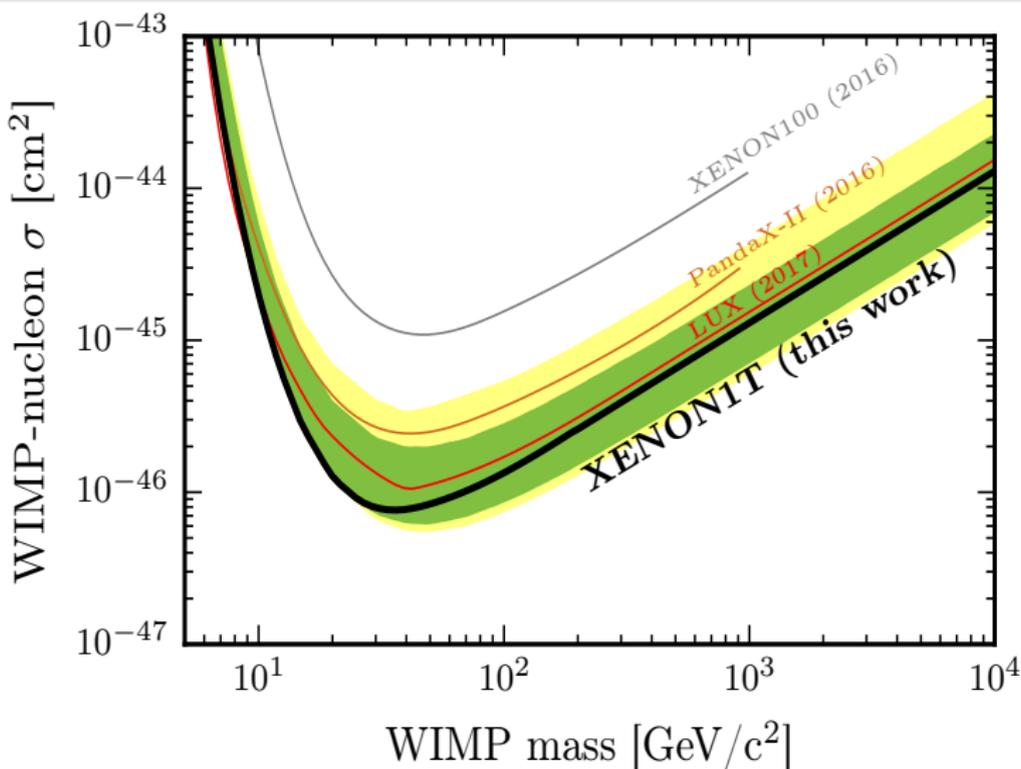


Figure : E. Aprile *et al.* [XENON Collaboration], Phys. Rev. Lett. **119**, no. 18, 181301 (2017), arXiv:1705.06655.

Where are all the WIMPs?

Direct detection vanishing mechanism: The Abelian case

Extension of the SM scalar sector by a scalar S :³

- Singlet under the G_{SM} .
- Charged under a global, softly broken $U(1)$ symmetry.

Invariant potential:

$$V_0 = -\frac{\mu_H^2}{2}|H|^2 + \frac{\lambda_H^2}{2}|H|^4 - \frac{\mu_S^2}{2}|S|^2 + \frac{\lambda_S}{2}|S|^4 + \lambda_{HS}|H|^2|S|^2.$$

Softly breaking term:

$$V_{\text{soft}} = -\frac{\mu'_S{}^2}{4}S^2 + \text{h.c.}$$

Observation:

There is one phase that can be absorbed in $S \rightarrow CP\text{-invariance!}$

³C. Gross, O. Lebedev and T. Toma, Phys. Rev. Lett. **119** (2017) no.19, 191801
doi:10.1103/PhysRevLett.119.191801 arXiv:1708.02253 [hep-ph].

The scalars develop VEVs:

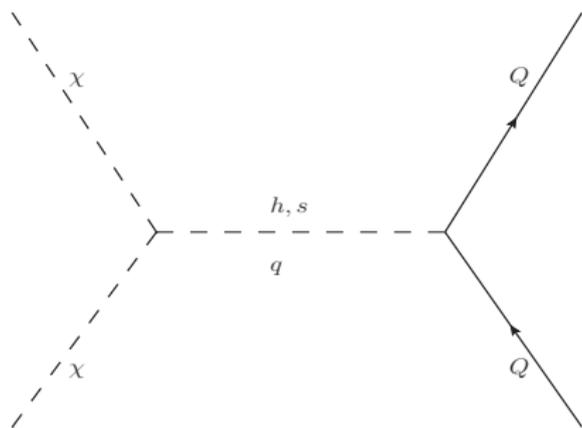
$$S = \frac{1}{\sqrt{2}}(v_S + s + i\chi),$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + h \end{pmatrix}.$$

Observation:

CP-invariance $\rightarrow \chi$ is stable!

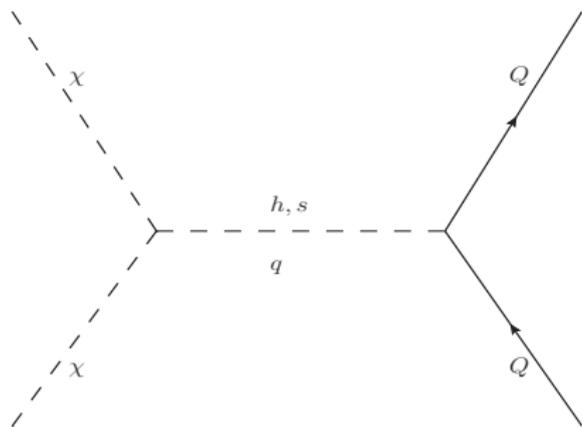
The cancellation mechanism



Propagator (at $q^2 \approx 0$) is proportional to the inverse of the CP-even mass matrix. So, the direct detection matrix element (A_{DD}) becomes

$$A_{DD} \sim (\lambda_{HS} v, \lambda_S v_S) \begin{pmatrix} \lambda_S v_S^2 & -\lambda_{HS} v v_S \\ -\lambda_{HS} v v_S & \lambda_{HV}^2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0.$$

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$$A_{DD} \sim (\lambda_{HS\nu}, \lambda_{S\nu S}) \begin{pmatrix} \lambda_{S\nu S}^2 & -\lambda_{HS\nu\nu S} \\ -\lambda_{HS\nu\nu S} & \lambda_{H\nu^2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = 0.$$

Goldstone bosons couple proportionally to the momentum (but some cancellations still have to occur).

The minimal model:

- Natural explanation for the missing DM signal.
- Natural emergence of a WIMP stabilizing symmetry.
- The origin of the softly broken term can be explained (gauged $U(1)$).
- WIMP at low masses ($\sim 10 - 10^4$ GeV).

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Generalization?

A simple Generalization

$U(1) \times S_N$

- N scalars (S_i).
- Transforming as $S_i \rightarrow e^{-ia} S_i$.
- Symmetry: $U(1) \times S_N$.
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This results to a pseudo-Goldstone:

$$\xi = \frac{1}{\sqrt{N}} \sum_{i=1}^N \chi_i \rightarrow S_N\text{-symmetric state.}$$

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Large number of phases \rightarrow no naturally stable particle.

N	#phases
1	1
2	3
≥ 3	$3 + \frac{1}{2}N(N - 1)$

Non-Abelian Generalization

A more attractive/simple case:

- One scalar, Φ .
- Doublet under a softly broken $SU(2)$.
- The potential is similar to the minimal $U(1)$ case (*i.e.* very simple).

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$$V_{\text{soft}} = \sum_{i=1}^2 \sum_{j=1}^2 \left[(m_{\Phi ij}^2 \Phi_i \Phi_j + \text{h.c.}) + m'_{\Phi ij}{}^2 \Phi_i^\dagger \Phi_j \right],$$

where $SU(2)$ is restored for $m_{\Phi ij} = m'_{\Phi ij} = 0$ and $m'_{\Phi ii} = \frac{\mu_{\Phi}}{2}$.

The doublet Φ acquires a VEV:

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi + i s \\ \rho + i \chi + v_\Phi \end{pmatrix},$$

resulting to mixing between the ρ and h .

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DM content:

- Three pseudo-Goldstone bosons.
- No CP-invariance (as in the minimal case), but emergence of a Z_2 .
- The interaction terms are invariant under orthogonal transformations \rightarrow All pseudo-Goldstone bosons are stable ($Z_2^{(1)} \times Z_2^{(2)} \times Z_2^{(3)}$).
- The $h - \rho$ mass matrix is similar to the minimal $U(1)$ case \rightarrow vanishing of the direct detection cross section for all DM particles!

The results of the $SU(2)$ case hold also for the $SU(N)$, at least for Φ in the fundamental representation.

$SU(N)$ generalization:

- There are $2N - 1$ pseudo-Goldstone bosons.
- The discrete symmetry is $Z_2^{(1)} \times Z_2^{(2)} \times Z_2^{(3)} \times \dots \times Z_2^{(2N-1)}$ ($2N - 1$ DM particles).

- There are models that can explain the missing WIMP signal.
- Naturally stable DM particles.
- It seems to be fairly easy to find such models (WIMP paradigm is still alive).
- No fine tuning.

- What happens if Φ is in another representation of $SU(2)$ (e.g. $SU(2)$ -triplet)?
- Phenomenological analysis (e.g. LHC detectability).
- Loop corrections to the scalar potential (also important for direct detection).
- The origin of the soft breaking terms.
- Any other models with the desirable attributes?

Thank you!

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