

Search for a heavy pseudoscalar boson A decaying into a Z and h boson in the semileptonic final state

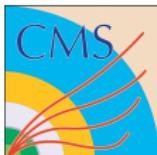
HEP 2018

"Recent Developments in High Energy Physics and Cosmology"

Emilios Ioannou¹, Alberto Zucchetta², Jehad Mousa¹, Panos Razis¹

¹University of Cyprus, ²University of Zurich

Athens, Greece, 28/3 -1/4/2018



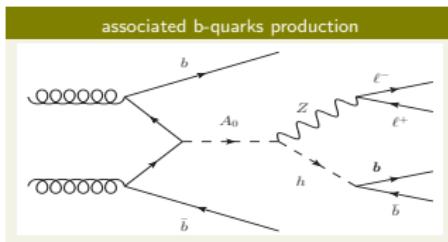
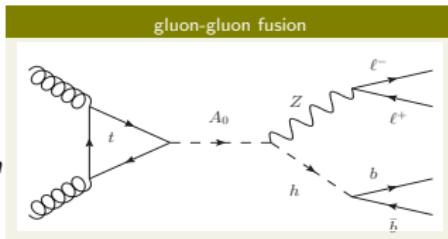
Universität
Zürich UZH



University
of Cyprus

Physics Motivations

- A pseudoscalar boson A is predicted in many BSM theories (2HDM, MSSM).
- In 2HDM (MSSM included), if $m_h + m_Z \lesssim m_A \lesssim 2m_t$, the branching ratio $A \rightarrow Z h$ might be important \rightarrow strong dependence on model parameters.
- Clean final state thanks to Z decays, large $\text{BR}(h \rightarrow b\bar{b})$.
- The SM-like Higgs boson is interpreted as the light scalar h with $m_h = 125 \text{ GeV}$.
- A pseudoscalar A could be generated by two Higgs mechanisms.



In this presentation

- ✓ Signal generation.
- ✓ Analysis strategy.
- ✓ Background estimation.
- ✓ Discriminators (angular, event shape).
- ✓ Systematic uncertainties.
- ✓ Results.

★ Results of blind analysis are being presented.

- Next-to-leading-order generators are available aMC@NLO.
 - 2HDMtII_NLO UFO model with QCD corrections for the type II 2HDM, both the top and bottom are massive and their yukawa are renormalized like their masses, i.e. on-shell.
- Possibility to fully simulate loop-induced processes (full top, bottom loop).
- Use **MadSpin** to have the correct $A \rightarrow Zh$, $Z \rightarrow \ell\ell$, $h \rightarrow b\bar{b}$ spin correlations.
- Mass points (500K events for each point):
 $m_A = 225, 250, 275, 300, 325, 350, 400, 500, 600, 700, 800, 900, 1000 \text{ GeV}$.
- Samples were generated for the processes:
 - $A \rightarrow Zh \rightarrow \ell\ell b\bar{b}$
 - $A \rightarrow Zh \rightarrow \nu\bar{\nu} b\bar{b}$
 - $b\bar{b}A \rightarrow Zh \rightarrow \ell\ell 2b2\bar{b}$
 - $b\bar{b}A \rightarrow Zh \rightarrow \nu\bar{\nu} 2b2\bar{b}$

Data and Monte Carlo samples

- Backgrounds are represented by all the processes with large MET or 2 leptons, and b-jets in the final state.

Data

SingleMu/Run2016(B-H)-23Sep2016-v*
SingleElectron/Run2016(B-H)-23Sep2016-v*
MET/Run2016(B-H)-23Sep2016-v*

Monte Carlo backgrounds

DYJetsToLL_M-50_TuneCUETP8M1_13TeV-amcatnloFXFX-pythia8
DYJetsToLL_M-50_HT-_TuneCUETP8M1_13TeV-madgraphMLM-pythia8
DY*JetsToLL_M-50_TuneCUETP8M1_13TeV-madgraphMLM-pythia8
DYJetsToNuNu_M-50_HT-_TuneCUETP8M1_13TeV-madgraphMLM-pythia8
WJetsToLNu_M-50_HT-_TuneCUETP8M1_13TeV-madgraphMLM-pythia8
TT_TuneCUETP8M1_13TeV-powheg-pythia8
TTToSemilepton_TuneCUETP8M1_13TeV-powheg-pythia8
TTTo2L2Nu_TuneCUETP8M1_13TeV-powheg-pythia8
ST_*_13TeV-amcatnlo-pythia8_TuneCUETP8M1
WWTo*_13TeV_amcatnloFXFX_madspin_pythia8
WZTo*_13TeV_amcatnloFXFX_madspin_pythia8
ZZTo*_13TeV_amcatnloFXFX_madspin_pythia8
ZH_HToBB_ZTo*_M125_13TeV_amcatnloFXFX_madspin_pythia8
WH_HToBB_WTo*_M125_13TeV_amcatnloFXFX_madspin_pythia8

Monte Carlo signal

GluGluToAToZhToLLBB_M*_13TeV-madgraphMLM-pythia8
GluGluToAToZhToNuNuBB_M*_13TeV-madgraphMLM-pythia8
BBAToZhToLLBB_M*_13TeV-madgraphMLM-pythia8
BBAToZhToNuNuBB_M*_13TeV-madgraphMLM-pythia8

- Golden JSON (35.9 fb^{-1}): Cert_271036-284044_13TeV_23Sep2016ReReco_Collisions16_JSON.txt
- All samples are from the Summer16 official CMS production
- Pile-up reweighting using the official recipe ($\sigma_{in} = 69200 \mu\text{b}$)
- LO samples corrected for NLO QCD and NLO EWK contributions

Analysis strategy

The analysis in a nutshell

- Identify and select the decay products in the detector:
 - Two opposite sign, same flavour, isolated leptons or large \cancel{E}_T .
 - Two jets tagged as originated from b-quarks.
- Build intermediate candidates Z ($\ell^+\ell^-$ or $\nu\bar{\nu}$) and h ($b\bar{b}$) up to the A candidate.
- The signal would manifest itself as a **narrow peak** in the mass distribution over the SM background continuum.
- Define appropriate **control regions** (CR) to check background shape and normalization with data by inverting selections.

Pre-selections of $Z \rightarrow \ell^+\ell^-$

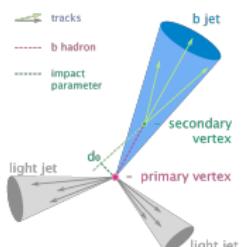
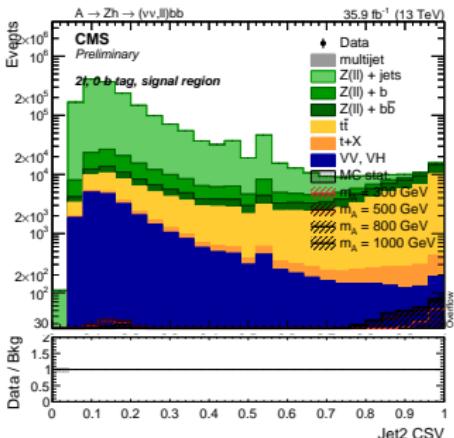
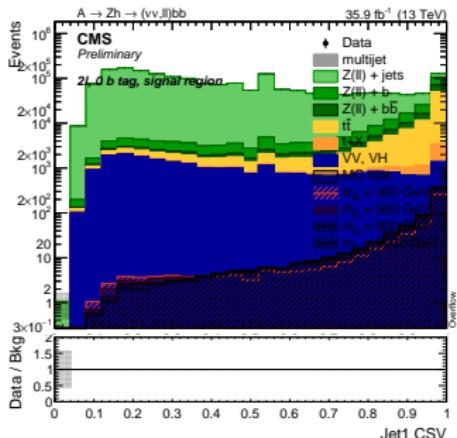
- At least two opposite sign, same flavour e or μ .
- At least two jets with $p_T > 30$ GeV.
- The Higgs candidate (h) is reconstructed from the highest-CSV jets.
- No boost requirements.

Pre-selections of $Z \rightarrow \nu\bar{\nu}$

- No isolated leptons (e, μ , τ).
- $\cancel{E}_T > 200$ GeV.
- At least two jets with $p_T > 30$ GeV.
- The Higgs candidate (h) is reconstructed from highest-CSV jets.
- $p_T^h > 200$ GeV (back-to-back).
- No jet closer in $\Delta\phi < 0.4$ to the \cancel{E}_T .

Combined-Secondary Vertex (CSV)

- The SM Higgs candidate will be reconstructed from the highest-CSV jets.
 - CSV: algorithm which identifies jets as b-jets, taking into account:
 - ✓ Kinematical and topological secondary vertex variables.
 - ✓ Information from track impact parameters.



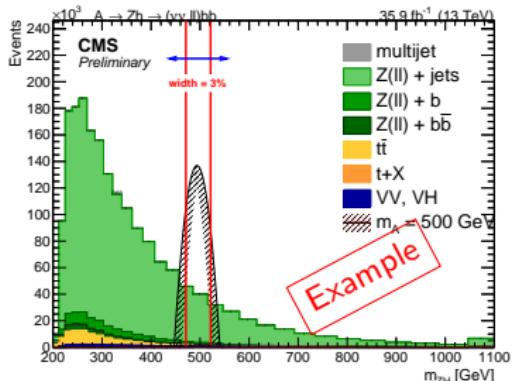
CSV working point selections

Working point	Cut
Loose	0.5426
Medium	0.8484
Tight	0.9535

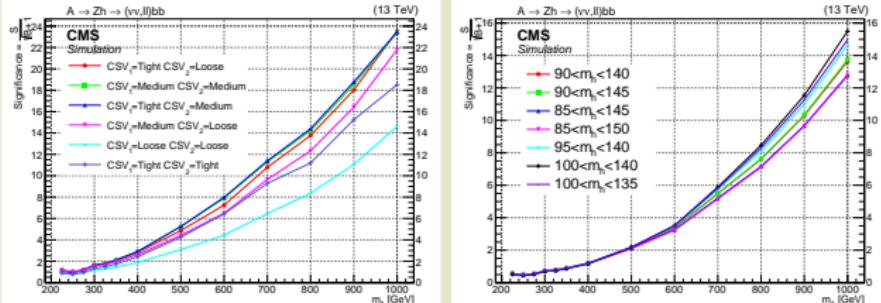
- A proper CSV working point selection should be defined.
- Which should be the most suitable criterion?
 - ✓ CSV working point selection → larger significance.

■ How do we compute the Significance?

(If one selection can reject background events from the signal region, while signal events remain unaffected)



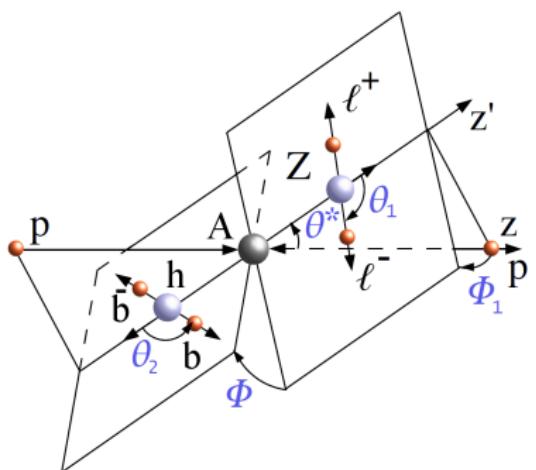
- ✓ Choosing a window with 3% width for each A candidate.
- ✓ Integrating the bins (signal and background) which are included in the window, separately.
- ✓ Using the formula $FOM = \frac{Signal}{\sqrt{Background+1}}$ to extract the **Significance** for each mass point.



- CSV working point selection:
(Medium, Medium) has been selected because of the larger significance.
- Higgs mass window:
(black line)
100 < m_h < 140 GeV → improved significance by up to 10-15%.

Angular discriminator

- In order to distinguish the signal from background, two discriminators are introduced:
 - Angular discriminator ($Z + b\bar{b}$).
 - Event shape discriminator ($t\bar{t}$).
- Angular Discriminator:
 - A distinctive signature of signal is that the Z boson is transversally polarized w.r.t. its flight direction for CP conservation
 - Additional discrimination comes from the $A \rightarrow Zh \rightarrow \ell\ell bb$ production ($\cos\theta^*$, Φ_1) and decay angles ($\cos\theta_1$, $\cos\theta_2$, Φ)



θ^* : angle between the A flight direction and the beam in the A rest frame.

θ_1 : angle between the ℓ^- and the Z flight direction in the Z rest frame.

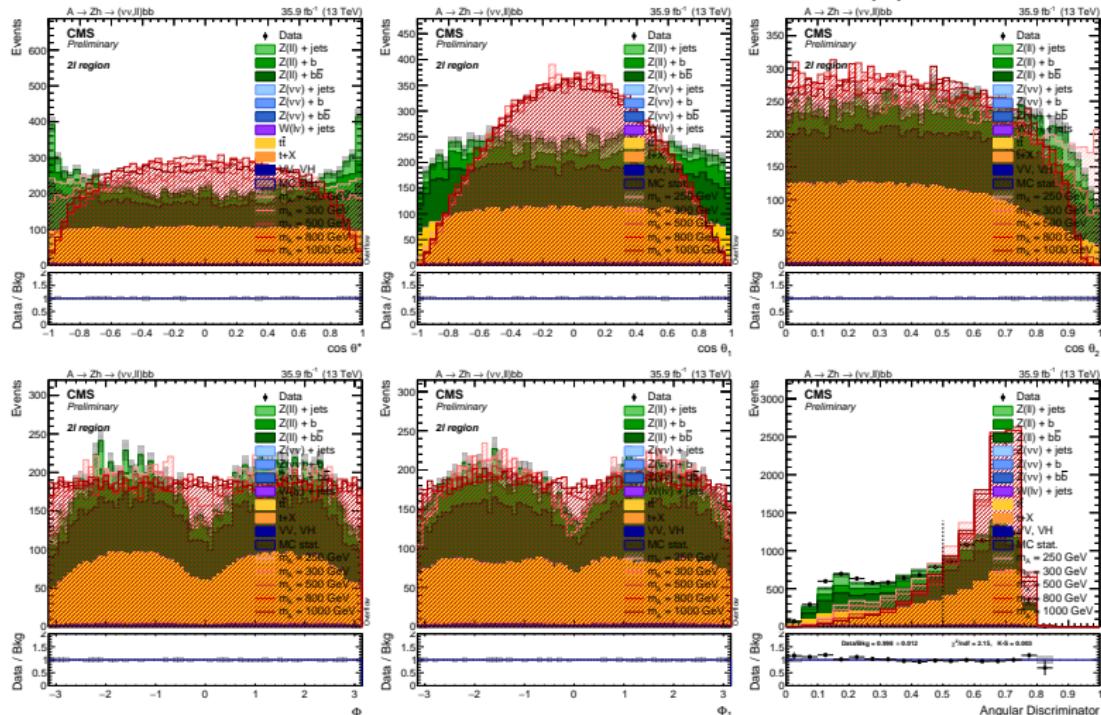
θ_2 : same as θ_1 , but for H and jets. The jet is chosen by convention.

Φ : angle between the Z and h decay planes.

Φ_1 : angle between the Z decay plane and the plane where the Z and beam directions lie.

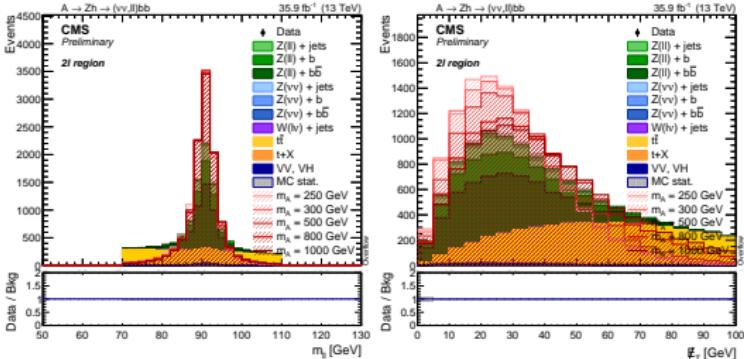
Angular variables

- $\cos \theta^*$ and $\cos \theta_1$ most discriminating, take advantage of the transversal polarization of the Z
- Likelihood discriminator trained against signal and Z+b(b) background.



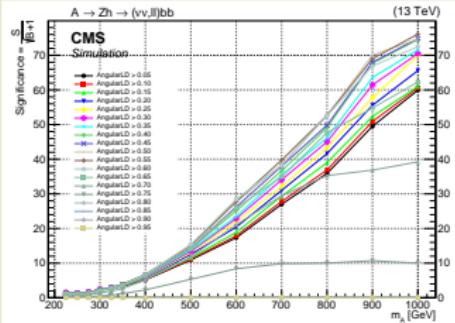
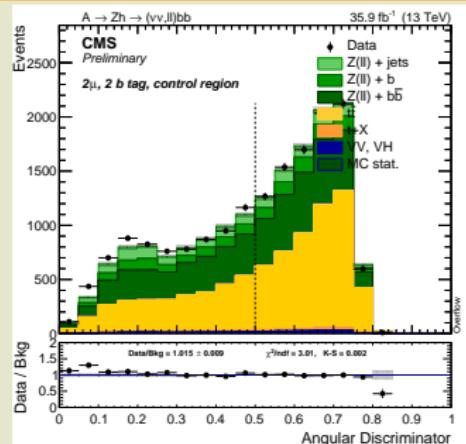
Event shape Discriminator

- Also $t\bar{t}$ is large ($1/3 \sim 1/2$) of the total background
- Dedicated discriminator trained against $t\bar{t}$
- Discriminating variables: Z mass, MET



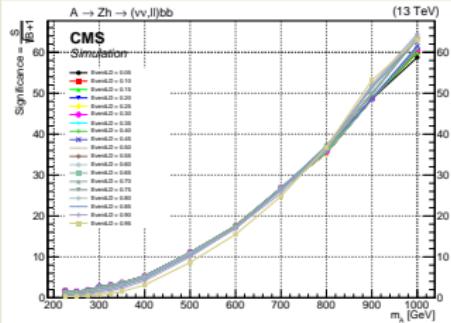
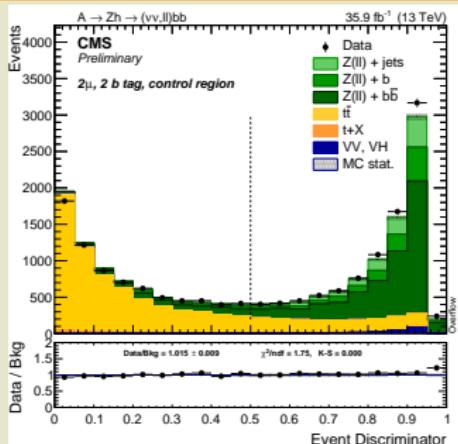
- Check the correlation matrix.
- Likelihood discriminator ensures optimal performance in absence of correlations between the variables.
- LD inputs also uncorrelated w.r.t. A boson mass.
- This is necessary to avoid any background sculpting when cutting on the discriminator.
- Correlation is very small → allowing implementation of cut on the discriminators.

Angular discriminator



■ Optimized selection: $\text{AngularLD} > 0.5$

Event shape discriminator

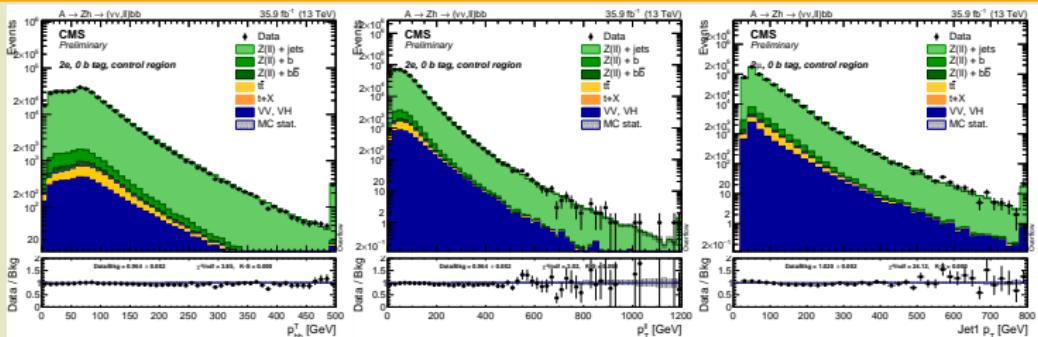


■ Optimized selection: $\text{EventLD} > 0.5$

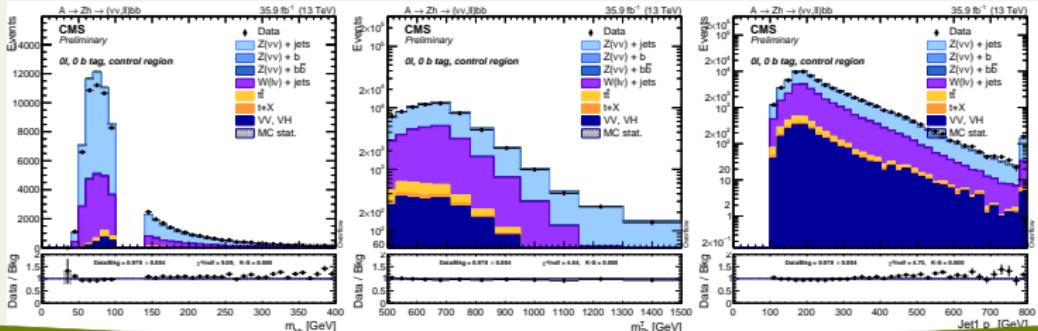
Data-driven background normalization (Z+jets)

Region	Z mass [GeV]	h mass [GeV]	CSV ₁	CSV ₂	\cancel{E}_T [GeV]
Z+jets	$70 < m_{\ell\ell} < 100$ e, μ, τ veto	$m_{bb} < 100, m_{bb} > 140$ $m_{bb} < 100, m_{bb} > 140$!Medium	!Medium	< 100
			!Medium	!Medium	> 200

$Z \rightarrow \ell^+ \ell^-$

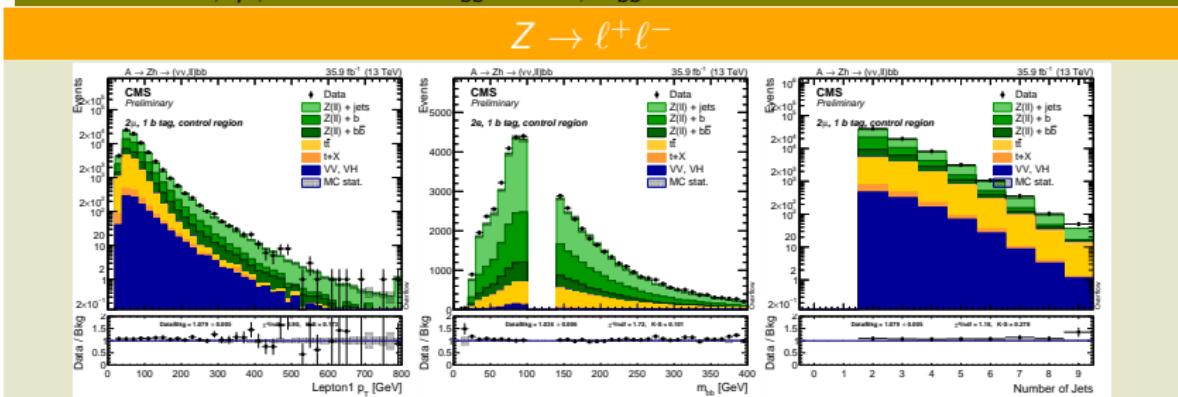


$Z \rightarrow \nu \bar{\nu}$

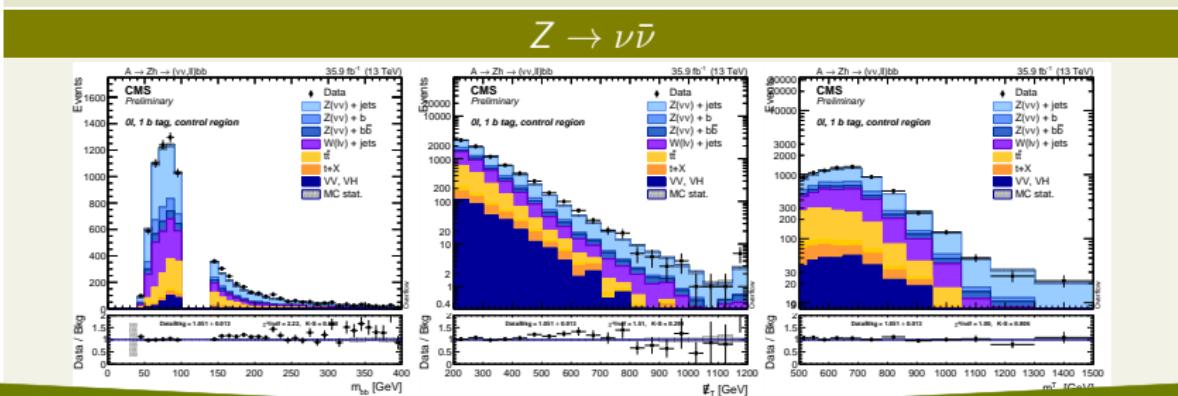


Data-driven background normalization ($Z+b$)

Region	Z mass [GeV]	h mass [GeV]	CSV_1	CSV_2	\cancel{E}_T [GeV]
$Z+b$	$70 < m_{\ell\ell} < 100$ e, μ, τ veto	$m_{bb} < 100, m_{bb} > 140$ $m_{bb} < 100, m_{bb} > 140$	Medium	!Medium	< 100 > 200
					$Z \rightarrow \ell^+\ell^-$



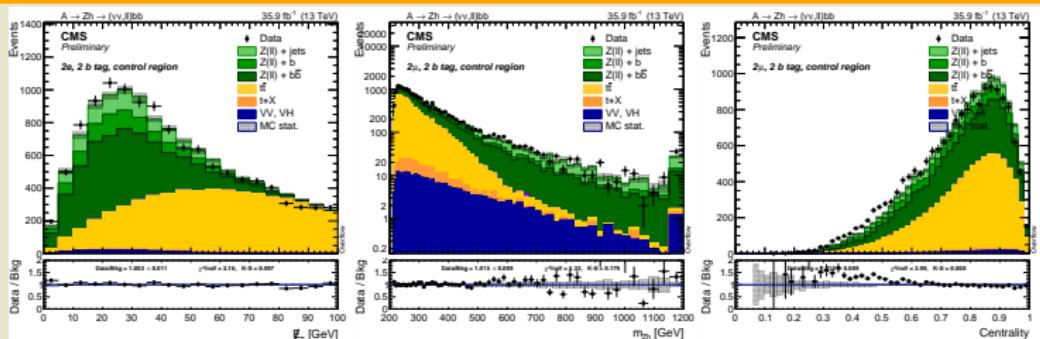
$Z \rightarrow \nu\bar{\nu}$



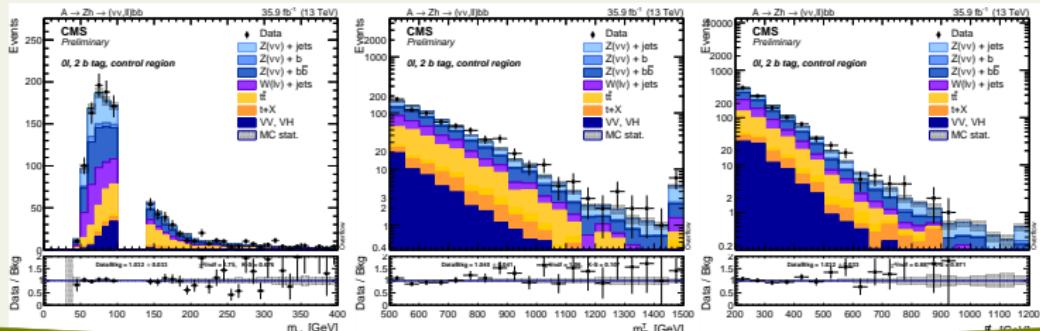
Data-driven background normalization ($Z + b\bar{b}$)

Region	Z mass [GeV]	h mass [GeV]	CSV ₁	CSV ₂	ℓ_T [GeV]
$Z + b\bar{b}$	$70 < m_{\ell\ell} < 100$	$m_{bb} < 100, m_{bb} > 140$	Medium	Medium	< 100
	0ℓ	$m_{bb} < 100, m_{bb} > 140$	Medium	Medium	> 200

$Z \rightarrow \ell^+\ell^-$



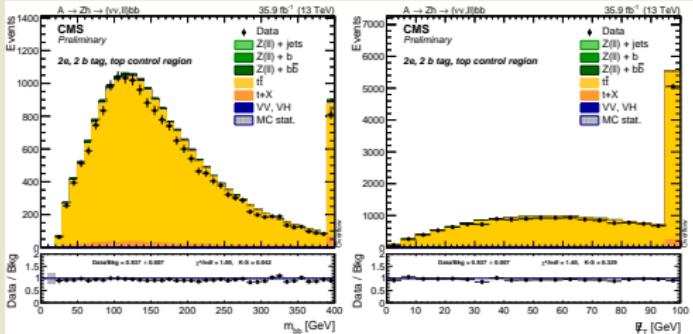
$Z \rightarrow \nu\bar{\nu}$



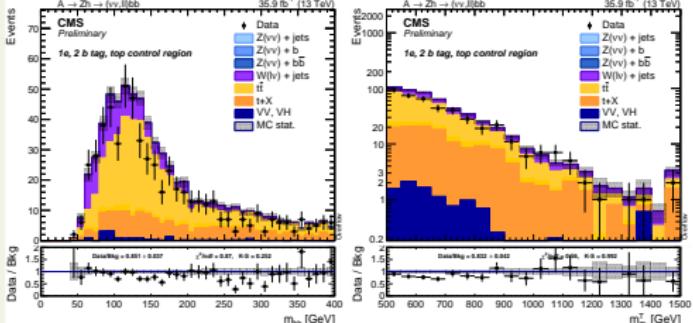
Data-driven background normalization (Top quark)

Region	Z mass [GeV]	h mass [GeV]	CSV ₁	CSV ₂	\cancel{E}_T [GeV]
Top and * W+jets	$m_{\ell\ell} < 70, m_{\ell\ell} > 110$ 1e or 1 μ	$m_{bb} < 100, m_{bb} > 140$ $m_{bb} < 100, m_{bb} > 140$	Medium	Medium	> 200

$Z \rightarrow \ell^+ \ell^-$



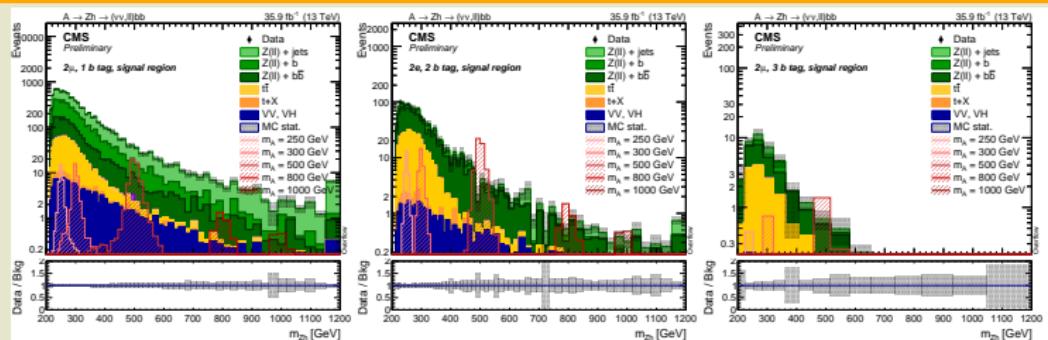
$Z \rightarrow \nu\bar{\nu}$



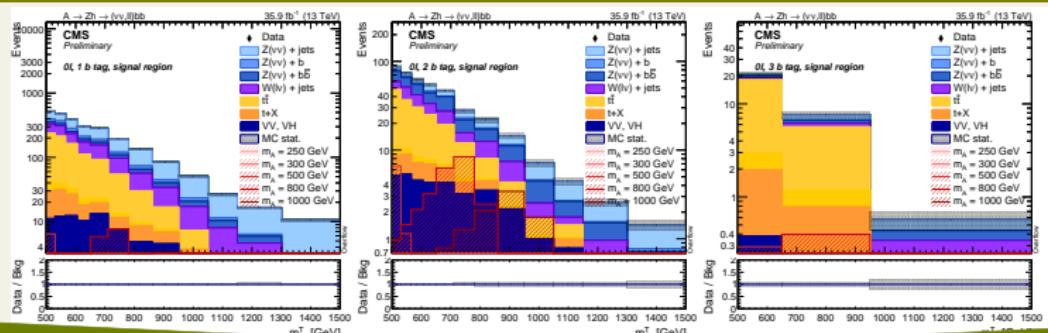
Signal Regions

Region	Z mass [GeV]	h mass [GeV]	CSV ₁	CSV ₂	\not{E}_T [GeV]
SR	$70 < m_{\ell\ell} < 100$	$100 < m_{bb} < 140$	Medium	Medium	< 100
	0ℓ	$100 < m_{bb} < 140$	Medium	Medium	> 200

$Z \rightarrow \ell^+ \ell^-$ (Three different signal regions)



$Z \rightarrow \nu \bar{\nu}$ (Three different signal regions)

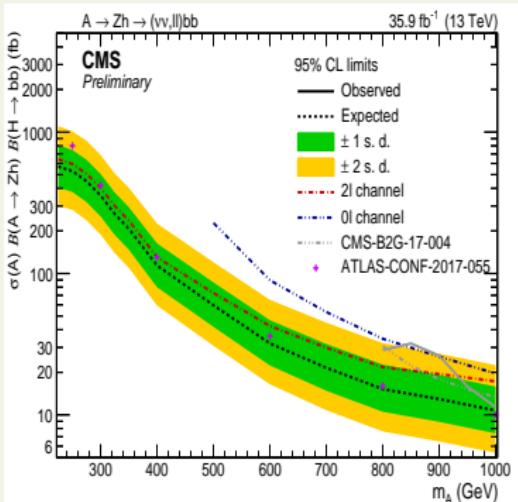


■ Summary of systematic shape uncertainties for backgrounds and signals

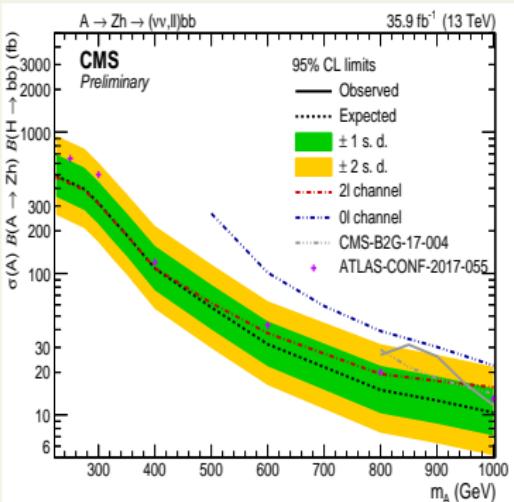
	Shape	Main backgrounds (Drell-Yan, $t\bar{t}$)	Other electroweak (single-top, VV, Vh)	Signal
Lepton and trigger efficiency	✓	-	2-3%	2-3%
Jet Energy Scale	✓	-	5%	2-6%
Jet Energy Resolution	✓	-	2%	1-2%
b-tagging	✓	-	4%	5-8%
Unclustered MET	✓	-	1%	1%
Pile-up	✓	-	1%	1%
PDF	✓	-	3-5%	4-8%
top quark p_T (only $t\bar{t}$)	✓	8-15%	-	-
Factorization and renormalization scale	✓		2-6%	6-14%
Monte Carlo modeling	✓		1-15 %	-
Monte Carlo statistics	✓		1-20%	-
Interpolation to SR			1-12%	-
Extrapolation to 3 b-tag SR		20-26% (3 b-tag only)		-
Cross section	-		2-10%	-
Luminosity	-		2.5%	2.5%

Results

gluon-gluon fusion



b-quark associated



- Expected limits with all systematics, included MC statistics.
- Gained $\sim 20\%$ from discriminator optimization.
- Lost 5 – 15% due to MC stat. uncertainty (depends on mass point).

Conclusions

In this presentation

- An overview of the $A \rightarrow Zh \rightarrow \ell^+\ell^-(\nu\bar{\nu})b\bar{b}$ analysis has been presented.
- All results which have been presented, are from blind analysis.
- After optimization of selections, the improvement is significantly better.
- The analysis is stable and optimized.

Thanks for your attention!

BackUp

Physics objects

Electrons

- Trigger: HLT_Ele27_WP*_Gsf_v* OR
HLT_Ele115_CaloIdVT_GsfTrkIdT_v*
 - SF with T&P package
- $p_T^{1,2} > 30, 10 \text{ GeV}$, $|\eta| < 2.5$
- Cut-based Loose Id
- PFIso included in Id
 - SF from EGamma POG

Muons

- Trigger: HLT_IsoMu24_v OR
HLT_IsoTkMu24_v
 - SF for trigger from Muon POG
- $p_T^{1,2} > 30, 10 \text{ GeV}$, $|\eta| < 2.4$
- At least one Tight, the other Loose
- PFIsolation $\Delta\beta$ corrected $< 0.15, 0.25$
 - SF for Id+Iso from Muon POG

Taus

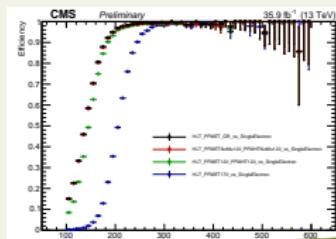
- $p_T > 18 \text{ GeV}$, $|\eta| < 2.5$
- Loose HPSPFTauId

Jets

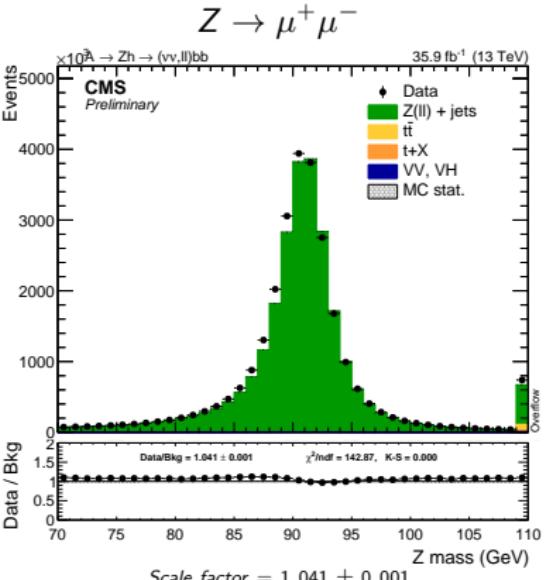
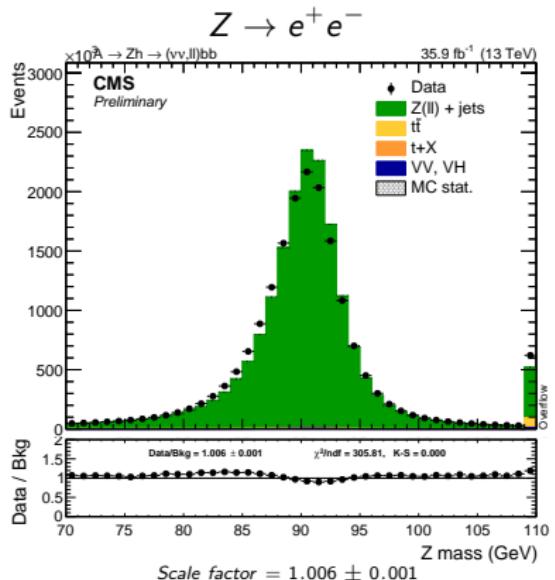
- AK4 PFJets with CHS
- default JEC L1L2L3(L2L3)
- $p_T > 30 \text{ GeV}$, $|\eta| < 2.5$
- PFLooseId
- $\Delta R_{\ell,j} > 0.4$
- CSVIVFV2 b-tagging
- SF from BTV POG

MET

- Trigger: HLT_PFMET110_PFMHT110_*
- 95% efficient for $MET > 200 \text{ GeV}$
- Type-I corrected MET
- Recommended filters applied

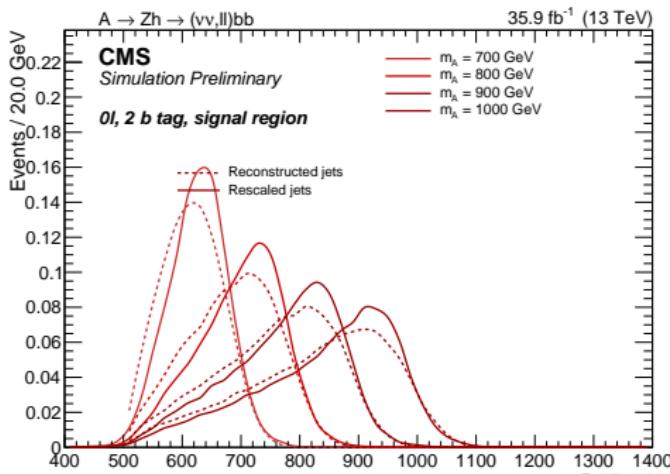
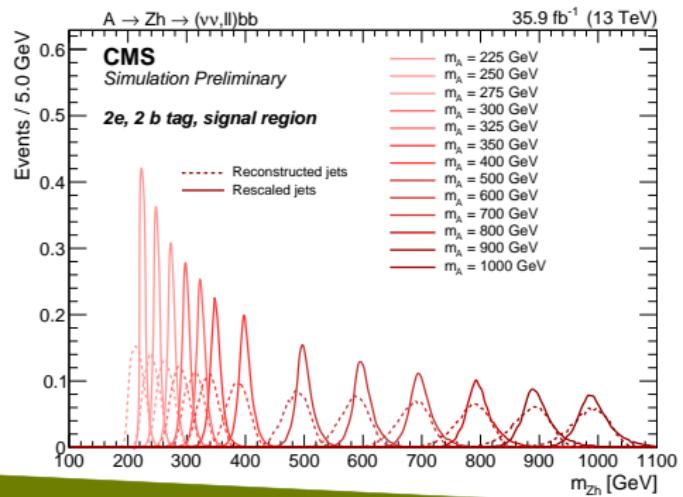


- Reconstruction, identification, isolation applied in both data and MC (with SF).
- Good out-of-the-box agreement: $\sim 4\%$ in $Z \rightarrow \mu^+ \mu^-$



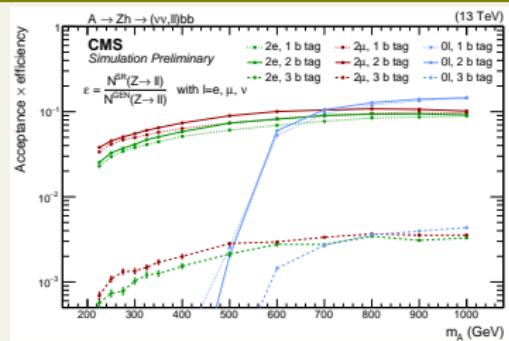
Kinematic constraints

- Exploit the measurement of the Higgs mass m_{bb} to rescale the jets' p_T to match $m_{jj} = 125$ GeV
- The rescale factor is proportional to the JEC uncertainty of each jet.
- New widths are about 1/4 for low masses, 1/2 for higher masses
 - at high m_A , the resolution is dominated by p_T^{bb}
- No so effective in the $Z \rightarrow \nu\nu$ channel, where the MET resolution is dominating
- Similar constrain is applied to the leptos' p_T to match m_Z , but the impact on the m_A resolution is much smaller.

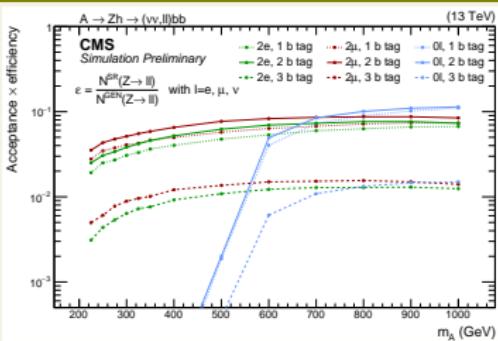


Signal Efficiency

gluon - gluon fusion



b-quarks associated



Comments

- Overall signal efficiency is 5-15%.
- Half of the events recovered in the 1 b-tag category.
- 3 b-tag category is useless for gluon-gluon fusion signal but is not negligible in the b-quark associated production ($\approx 20\%$ fraction of events with the b-quarks in acceptance).

Systematic uncertainties

Jet Energy Scale:

JES is varied by its uncertainty and propagated to composite objects (\cancel{E}_T). Effect on the A peak is negligible after the kinematic fit.

Jet Energy Resolution:

JER is evaluated by smearing the jet four-momentum according to the measured uncertainty.

Missing Transverse Energy:

We consider standard recipes for Unclustered MET only. Variations for JES, JER and leptons already included in the corresponding systematics.

Pile-up and PDF:

Followed standard prescriptions as reported in TWikis.

Systematic uncertainties (continue)

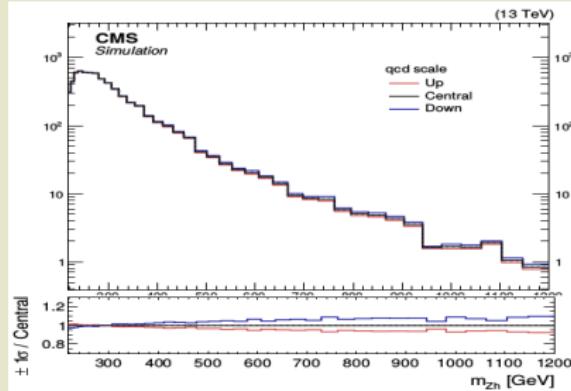
Lepton uncertainties:

propagated effect of Trigger, Reco, Id and Isolation scale factors.

	Trigger	Id + Iso (1 ℓ)	Id + Iso (2 ℓ)
Electrons	0.9% (2%)	1.9%	3.6%
Muons	3% (0.5%)	1.0% (2%+1.5%)	1.9% (4%)

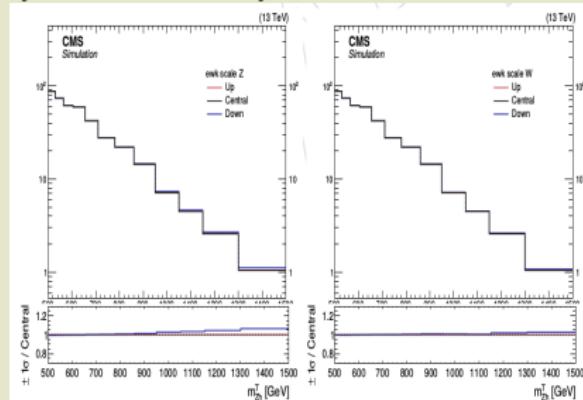
QCD renormalization and factorization scale:

MC reweighted charging Q by factor 2.



Electroweak NLO corrections:

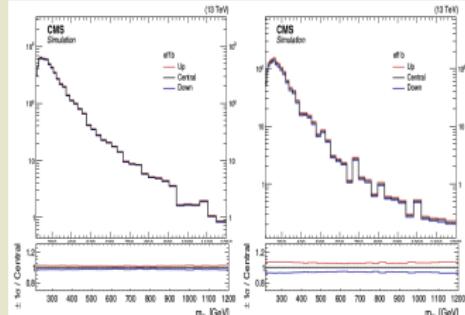
Take the full effect of the correction as systematic uncertainty.



Systematic uncertainties (continue)

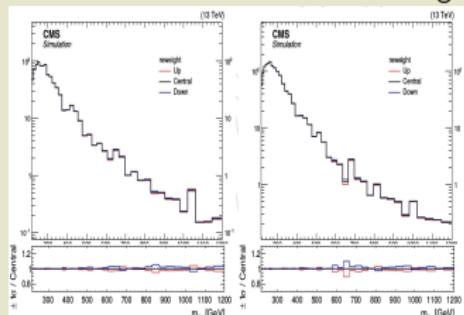
B-tagging:

Scale factors are varied by their error, weights propagated to the final distributions.



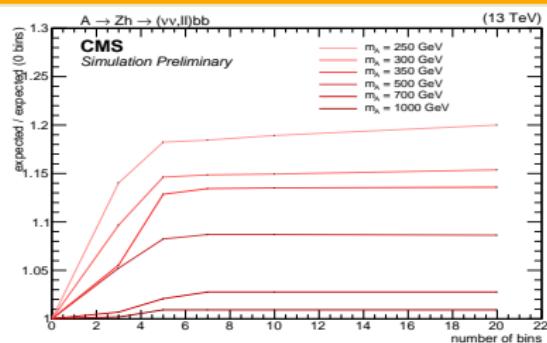
MC modeling:

Reweight in the Centrality variable to give the background a degree of freedom and cover minor mismodelings.



MC statistics:

Bins varied by their statistic uncertainty for each background independently.
 Only the 5 bins with the highest S/B are considered.



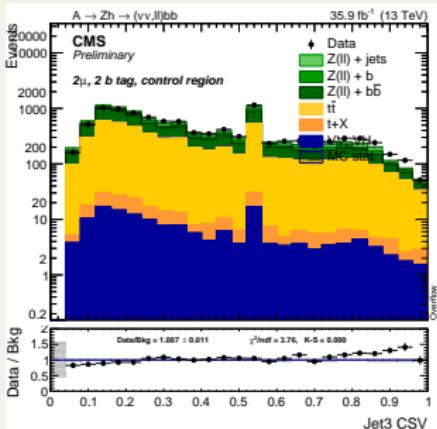
Interpolation to the SR

Interpolation to the SR

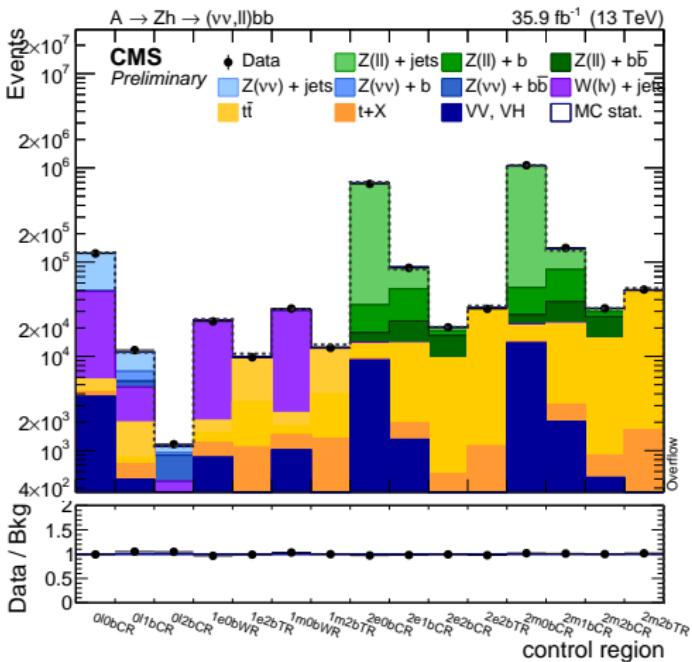
- In some CR, there is a non-negligible mismodeling in m_{bb}
- This may lead to a bias in the estimation of the background normalization
- Take the difference between the lower and higher sideband as systematic uncertainty

	1 b-tag	2 b-tag	3 b-tag
0ℓ	12.1%	1.2%	3.6%
$2e$	1.3%	5.4%	0.7%
2μ	1.2%	4.7%	7.1%

- Mismodeling in 3rd jet CSV observed consistently in all control regions
- Likely that this effect also yields an excess of data in the 3 b-tag SR
- Take the difference from 1 between the data and bkg simulation in 3 b-tag control regions (exactly as the 2 b-tag CR, but with 3 or more b-tags) as systematic uncertainty
- The difference is 20-25% in all final states



- A background-only fit is performed deriving the background scale factors.



2HDM in nutshell

- Two type of 2HDM are taken into account:

Type I: all quarks coupled to the same Higgs doublet (Φ_2). Discrete symmetry ($\Phi_1 = -\Phi_1$).

Type II: Right-handed quarks ($Q = 2/3$) coupled to Φ_2 . Right-handed quarks ($Q = -1/3$) and leptons coupled to Φ_1 .

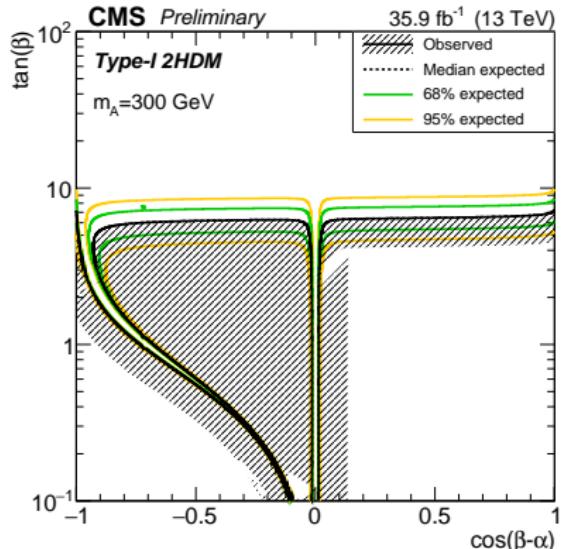
- Large number of free parameters: m_h , m_H , m_A , m_{H^\pm} , m_{12} , β , α , $\lambda_{6,7}$
- Recommendations from the Higgs Working Group:

- $m_H = m_A = m_{H^\pm}$: A searches not really dependent on m_H , m_{H^\pm} anyway
- $m_{12}^2 = m_A \frac{\tan \beta}{1 + \tan^2 \beta}$: discrete Z_2 symmetry broken as in MSSM
- $\lambda_{6,7} = 0$ to avoid CP-violation at tree level

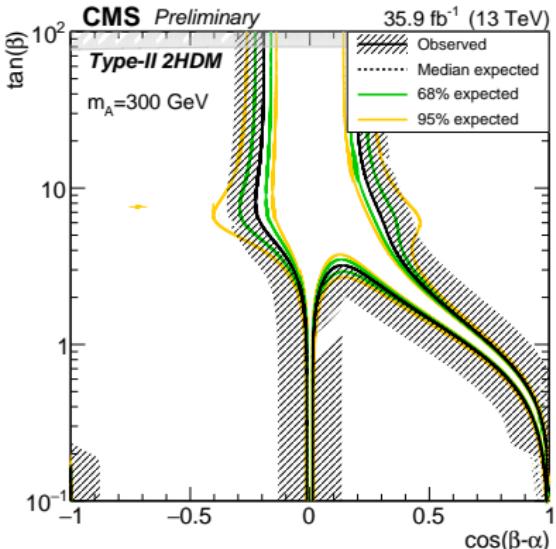
- Two parameters left other than m_A :
 - α : mixing angle of the two doublets
 - $\tan \beta$: the ratio of the vev of the two doublets

- For each mass point, perform a scan in $\tan \beta$ and $\cos(\beta - \alpha)$
 - $0.1 < \tan \beta < 100$,
 - $-1 \leq \cos(\beta - \alpha) \leq 1$ $0 \leq \beta - \alpha \leq \pi$
- Cross sections and Br calculated at NNLO with SuShI 1.6.1 and 2HDMC 1.7.0

Interpretation in 2HDM



Expected limit on $\sigma \times Br$ in Type I 2HDM



Expected limit on $\sigma \times Br$ in Type II 2HDM

- Expected exclusion up to $\tan \beta \approx 6 \sim 7$ in Type I.
- Large parameter space excluded in Type II, also at high $\tan \beta$.