

**FINITE QUANTUM MECHANICS
ON THE HORIZON OF BLACK HOLES**

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**CONFERENCE IN MEMORY OF
IOANNIS BAKAS**

RECENT WORK

- Modular discretization of the AdS₂/CFT₁ Holography
M.Axenides,E.Floratos,S.Nicolis
JHEP 1402(2014)109 arXiv:1306.5670
- Chaotic Information Processing by Extremal Black Holes
M.Axenides,E.G.Floratos,S.Nicolis
Int. J. Mod. Phys. D24 (2015) 1542.0122
arXiv:1504.00483

- Quantum cat map dynamics on AdS2
Minos Axenides, Emmanuel Floratos,
Stam Nicolis
- arXiv:1608.07845
- Arithmetic Circuits for Multilevel Qudits Based
on Quantum Fourier Transform
- Archimedes Pavlidis, Emmanuel Floratos
- . arXiv:1707.08834

ARITHMETIC GEOMETRY

A PROPER FRAMEWORK FOR THE BH QUANTUM INFORMATION PARADOX

- ASSUMPTIONS
- FINITENES OF BH ENTROPY $S \rightarrow$
FINITE DIMENSIONAL HILBERT SPACE OF BH MICROSCOPIC
STATES $\text{Dim}[H]=\text{Exp}[S]$
- FINITE DIMENSIONAL HILBERT SPACE \rightarrow DISCRETE AND
FINITE SPACE-TIME
- THE SIMPLEST MODEL FOR SINGLE PARTICLE DYNAMICS A
DISCRETE AND FINITE PHASE SPACE $Z[N] \times Z[N]$
 $Z[N]=\{0,1,2,3,\dots,N-1\}$ ALL INTEGERS MODULO N
FOR $N=\text{PRIME INTEGER}=p$ GALOIS FIELD $F[p]$

ARITHMETIC GEOMETRY FOR HORIZON BH DYNAMICS

A MODEL FOR PLANCK SCALE SPACE-TIME

- CONSEQUENCES
- MODULAR ARITHMETIC → RANDOM POINT GEOMETRIES
- MODULAR DYNAMICS $A = 2 \times 2$ MODULAR MATRICES IN $SL[2, \mathbb{Z}[N]]$
→ CHAOS AUTOMATIC
- FINITE DIMENSIONAL HILBERT SPACES
- FINITE QUANTUM MECHANICS AND QUANTUM FIELD THEORIES
- $N = p^n$, QUDITS OF INFORMATION,
- FINITE DIMENSIONAL UNITARY RANDOM EVOLUTION OPERATORS $U[A]$
- QUANTUM COMPLEXITY → COUNTING THE GATES OF
- QUANTUM CIRCUITS FOR CONSTRUCTING $U[A]$

- MOTIVATION

1) THE EIGENSTATE THERMALIZATION HYPOTHESIS
GAUSSIAN PDF OF EIGENSTATE'S PROB VALUES
FLAT PDF OF EIGENSTATE'S PHASES

2) SATURATION OF THE SCRAMBLING TIME BOUND

3) RELATION OF QUANTUM COMPLEXITY WITH
BH ENTROPY

SIMPLEST EXAMPLE

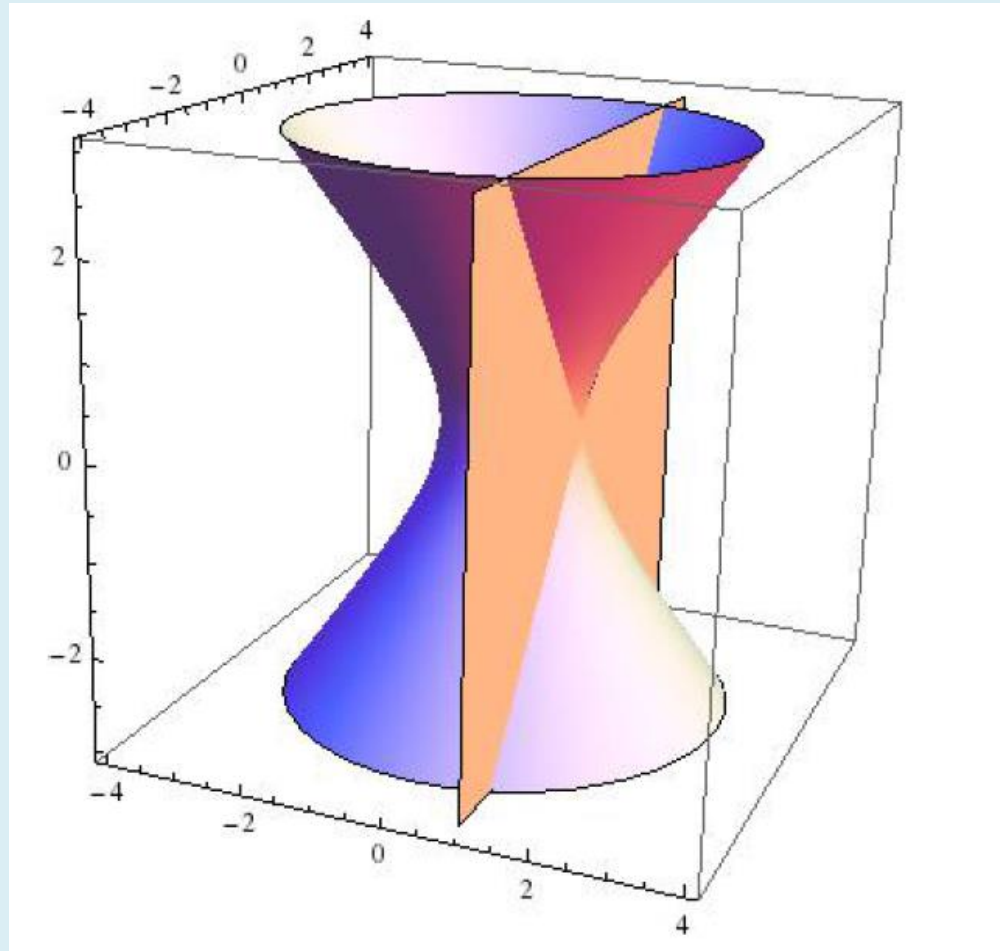
DIFUSION OF SINGLE PARTICLE WAVE PACKETS ON ADS2 TIME-RADIAL GEOMETRY OF EXTREMAL BH'S

- OLD AND LARGE NEAR EXTREMAL BH'S
- GEOMETRY = AdS2 X Σ , Σ =COMPACT ANGULAR DIRECTIONS
- ADS2 RADIAL MOTION
- AdS2[R]=SL[2,R]/SO[1,1,R]
- DISCRETIZE
- \square AdS2[N]=SL[2,ZN]/SO[1,1,ZN]
- CONSTRUCT THE AdS2 UNITARY EVOLUTION MATRIX OF PROBE STRING BITS
-> USE SUPERDCONFORMAL QM OF SL[2,ZN] ISOMETRY QUANTUM MAPS (TOWNSEND-STROMINGER-KALOSH 1998)

AdS2

NEAR HORIZON DEOMETRY OF EXTREMAL BH S

$$x_0^2 + x_1^2 - x_2^2 = 1$$



WEYL ACTION OF $SL(2, \mathbb{R})$ ON AdS_2

To every point $x_\mu \in AdS_2$, $\mu = 0, 1, 2$, we assign the traceless and real, 2×2 matrix

$$M(x) \equiv \begin{pmatrix} x_0 & x_1 + x_2 \\ x_1 - x_2 & -x_0 \end{pmatrix} \quad (2.3)$$

Its determinant is $\det M(x) = -x_0^2 - x_1^2 + x_2^2 = -1$.

The action of any $A \in SL(2, \mathbb{R})$ on AdS_2 is defined through the non-linear mapping

$$M(x') = AM(x)A^{-1} \quad (2.4)$$

This induces an $SO(1, 2)$ transformation on $(x_\mu)_{\mu=0,1,2}$,

$$x' \equiv L(A)x \quad (2.5)$$

Choosing as the origin of coordinates the base point $\mathbf{p} \equiv (1, 0, 0)$, its stability group $SO(1, 1)$ is the group of Lorentz transformations in the $x_0 = 0$ plane of $\mathcal{M}^{1,2}$ or equivalently, the “scaling” subgroup D of $SL(2, \mathbb{R})$

$$D \ni S(\lambda) \equiv \begin{pmatrix} \lambda & 0 \\ 0 & \lambda^{-1} \end{pmatrix} \quad (2.6)$$

for $\lambda \in \mathbb{R}^*$.

For this choice of the stability point, we define the coset h_A by decomposing A as

$$A = h_A S(\lambda_A) \quad (2.7)$$

Thus, we associate uniquely to every point $x \in AdS_2$ the corresponding coset representative $h_A(x)$.

**ARITHMETIC DISCRETIZATION OF $AdS_2=SL[2,R]/SO[1,1]$
 $\Rightarrow AdS_2[N]=SL[2,Z[N]]/SO[1,1,[Z[N]]]$**

$$X_0^2 + X_1^2 - X_2^2 = 1 \pmod{N}$$

ALL INTEGER SOLUTIONS $\pmod{N} \Rightarrow$ DISCRETE SET OF POINTS = $AdS_2[N]$

$$X_0 = A - M B,$$

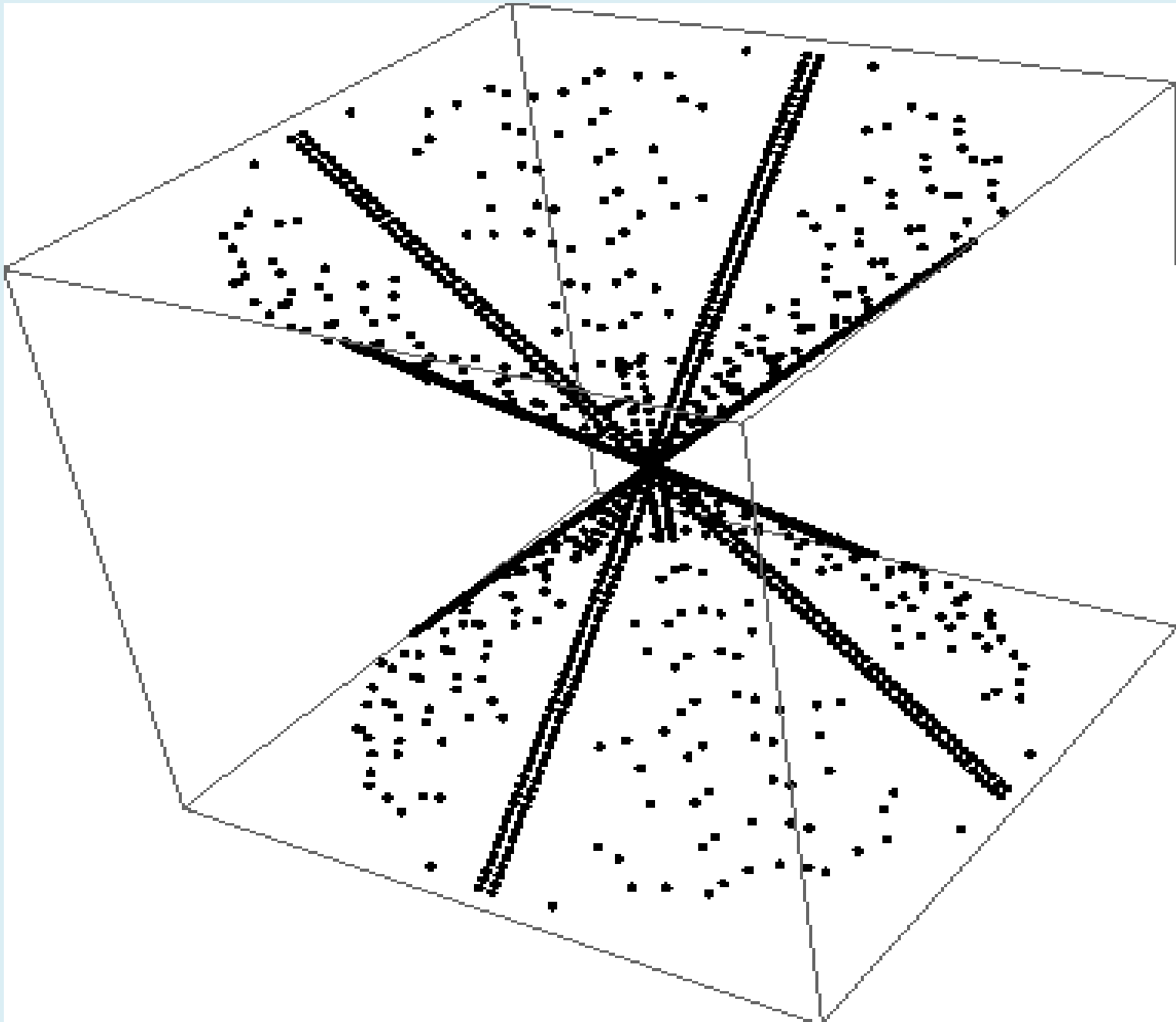
$$X_1 = B + M A$$

$$X_2 = M \quad A, B, M = 0, 1, \dots, N-1$$

$$A^2 + B^2 = 1 \pmod{N}, \quad \text{DISCRETE CIRCLE } S_1[N]$$

ROTATING THE 'LINE', $\{ X_0=0, X_1=M, X_2=M \mid M=0, 1, 2, \dots, N-1 \}$ AROUND $S_1[N]$

$SO[1,1,Z[N]]$ STABILITY GROUP OF $P: X_0=1, X_1=X_2=0$



DISCRETE SUPERCONFORMAL CONFORMAL DYNAMICS

ARNOLD CAT MAP $A = \{\{1,1\}, \{1,2\}\},$

WEYL ACTION ON $AdS_2[N]$ $X \rightarrow A.X.A^{-1}$

PROPERTIES

a) STRONG ARITHMETIC CHAOS

(ARNOLD, FORD, BERRY VOROS, VIVALDI, DI VIZENZO)

b) HOLOGRAPHY \Rightarrow NON LOCAL REDUNDANT STORAGE OF INFORMATION

c) MIXING TIME - LIAPUNOV EXPONENT

d) GENERATION OF KOLMOGOROV - SINAI ENTROPY

BASIC PROPERTY OF ARNOLD CAT MAP = FIBONACCI CHAOS

$$A^n = \{\{f(2n-1), f(2n)\}, \{f(2n), f(2n+1)\}\},$$

$f(n)$ FIBONACCI INTEGERS

$$f[n+1] = f[n] + f[n-1]; f[0] = 0, f[1] = 1$$

FOR ANY INTEGER N

Periods of A MOD[N] $A^{T(N)} = \text{IdentityMatrix MOD}[N]$

DYSON: IF $N = f[m] \rightarrow T(N) = 2m$

BUT SINCE FOR $m \rightarrow \text{INFINITY}$ $f[m] \rightarrow \text{Exp}[c m], c = \log[\text{Goldenratio}]$

WE OBTAIN $T[N] \rightarrow \text{Log}[N]$

LOGARITHMIC TIME CHAOTIC MIXING (SCRAMBLING)

For FIBONACCI SEQUENCE OF INTEGER HILBERT SPACE DIMENSIONS

**EIGENSTATE THERMALIZATION SENARIO
PAGE,DEUTSCH,BERRY,SREDNICKI**

**IF THE EIGENSTATES OF A CLOSED QM SYSTEM ARE RANDOM
(RANDOM PHASES AND GAUSSIAN DISTRIBUTED AMPLITUDES)**

**THEN ANY INITIAL PURE STATE OF A SUBSYSTEM THERMALIZES
TO
THE THERMAL DENSITY MATRIX OF THE SUBSYSTEM**

RELATION TO THE INFORMATION PARADOX

2013 SREDNICKI TALK TO KITP

ARNOLD QUANTUM CAT MAP $A=\{\{1,1\},\{2,1\}\}$

EXACT CONSTRUCTION OF THE SPECTRUM AND EIGENSTATES FOR $N=p$, prime,

LINEAR SPECTRUM

RANDOM EIGENSTATES \rightarrow LINEAR COMBINATIONS OF MULTIPLICATIVE CHARACTERS OF $GF[P]$

1) RANDOM PHASES

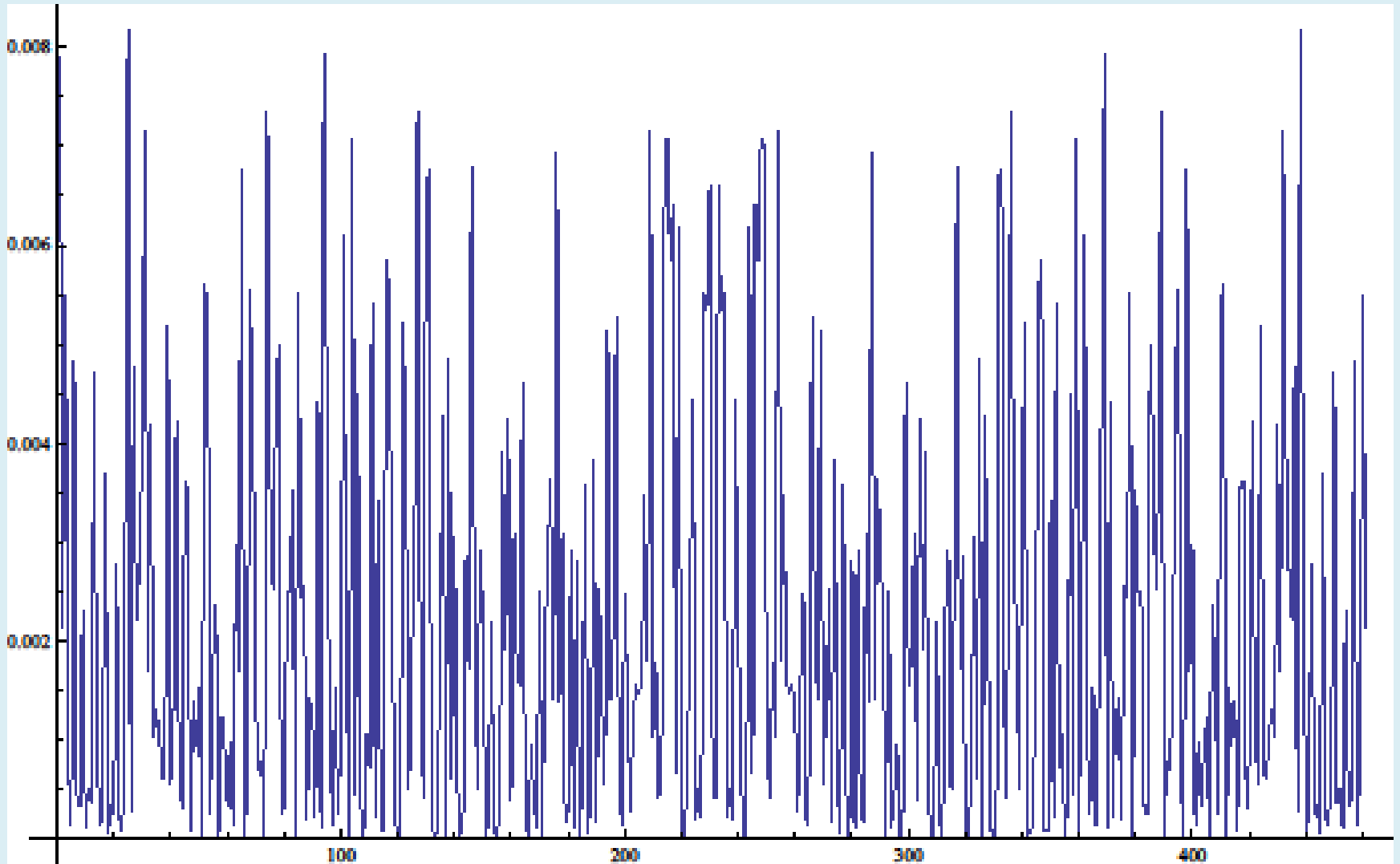
2) GAUSSIAN RANDOM AMPLITUDES

BUT SCARS (VOROS., NONEMACHER) F
FOR SEQUENCES OF N 's WITH SHORT PERIODS

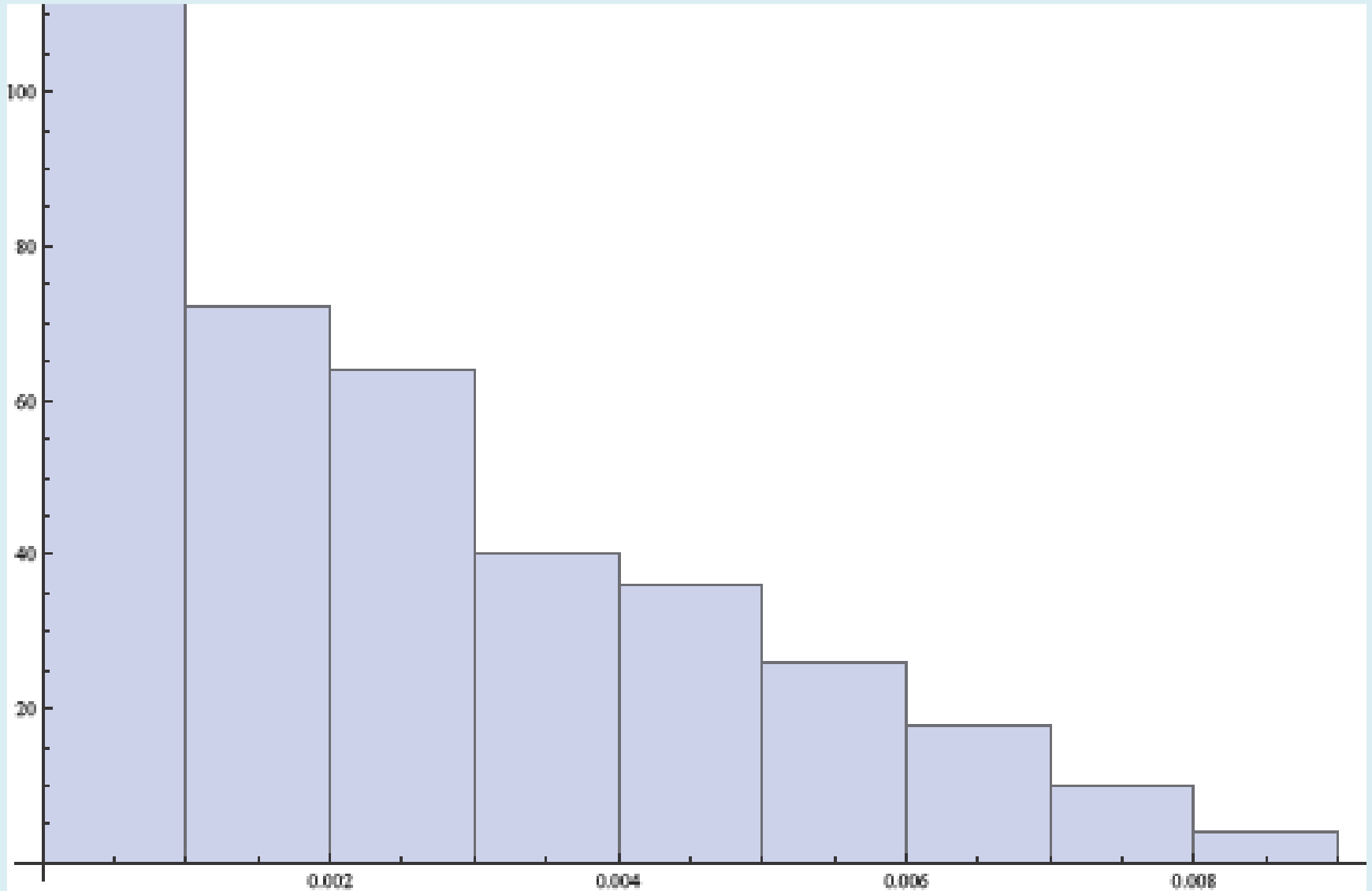
QUANTUM CHAOS, STRONG MIXING

- FACTORIZATION FOR ARNOLD CAT MAPS IMPLIES LOGARITHMIC IMPROVEMENT FROM $N^2 \rightarrow N \log N$
- USING QUANTUM CIRCUITS FOR THE IMPLEMENTATION OF THE QUANTUM MAP AND COUNTING THE NUMBER OF GATES $N \log N \rightarrow (\log N)^2$
- EXACTLY AS FOR THE QUANTUM FOURIER FACTORIZATION ALGORITHM OF SHOR.

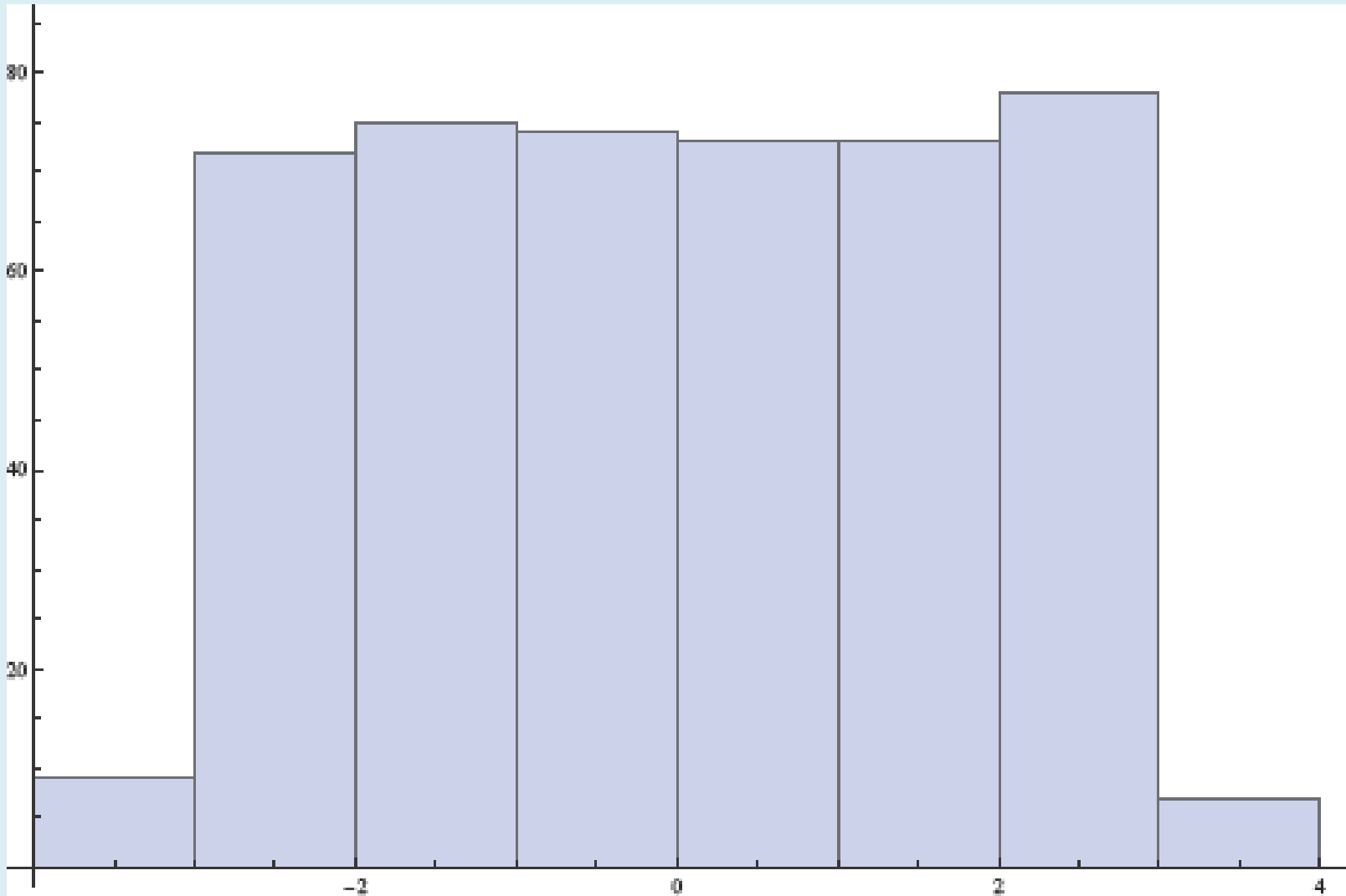
- $N=461, T[461]=23$,GROUND STATE, $\Phi=0$



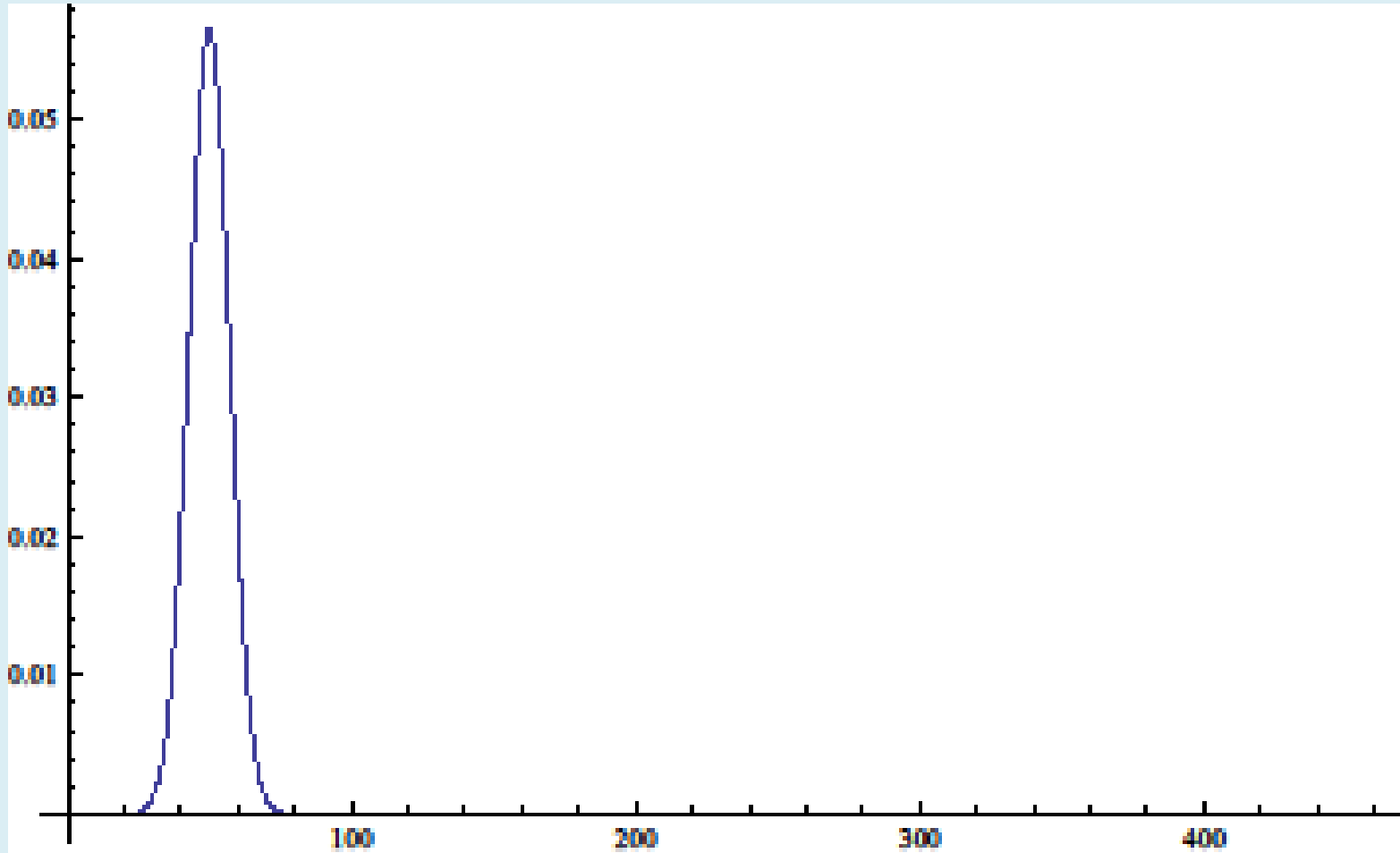
- GROUND STATE AMPLSQUARE DF



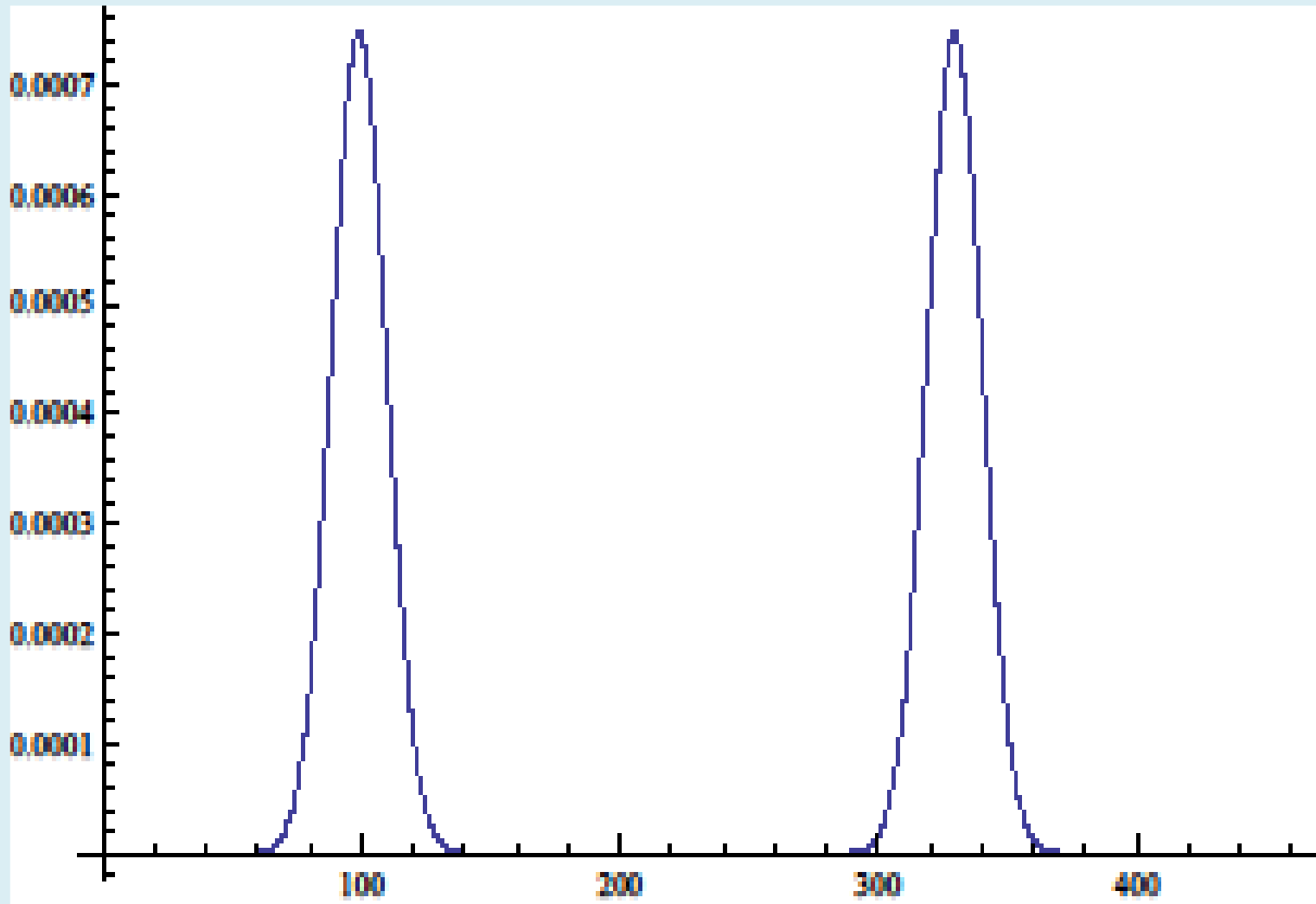
- GROUND STATE PHASE DF



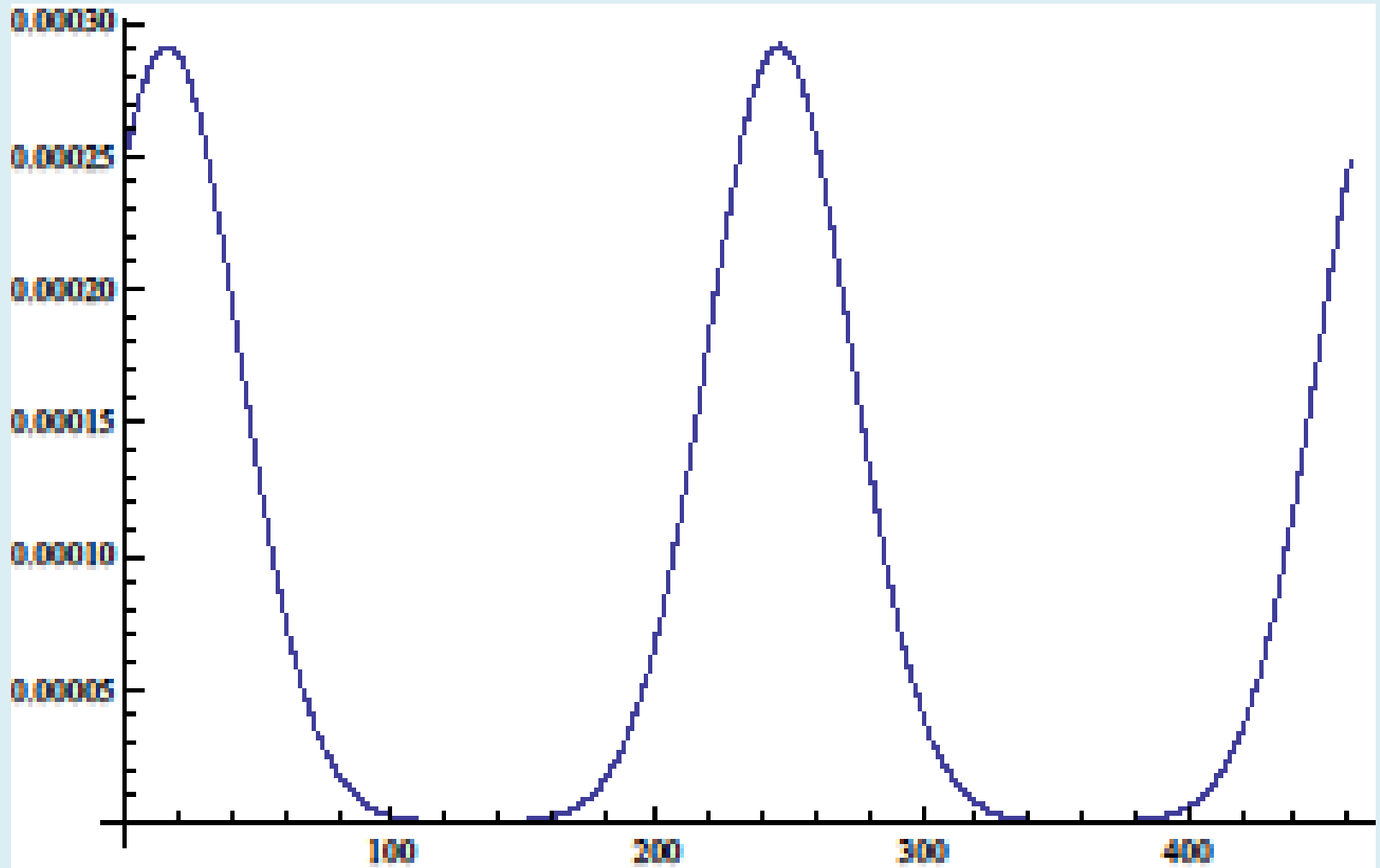
- SCATTERING EXPERIMENT, $N=p=461$
- GAUSSIAN WF AT $T=0$



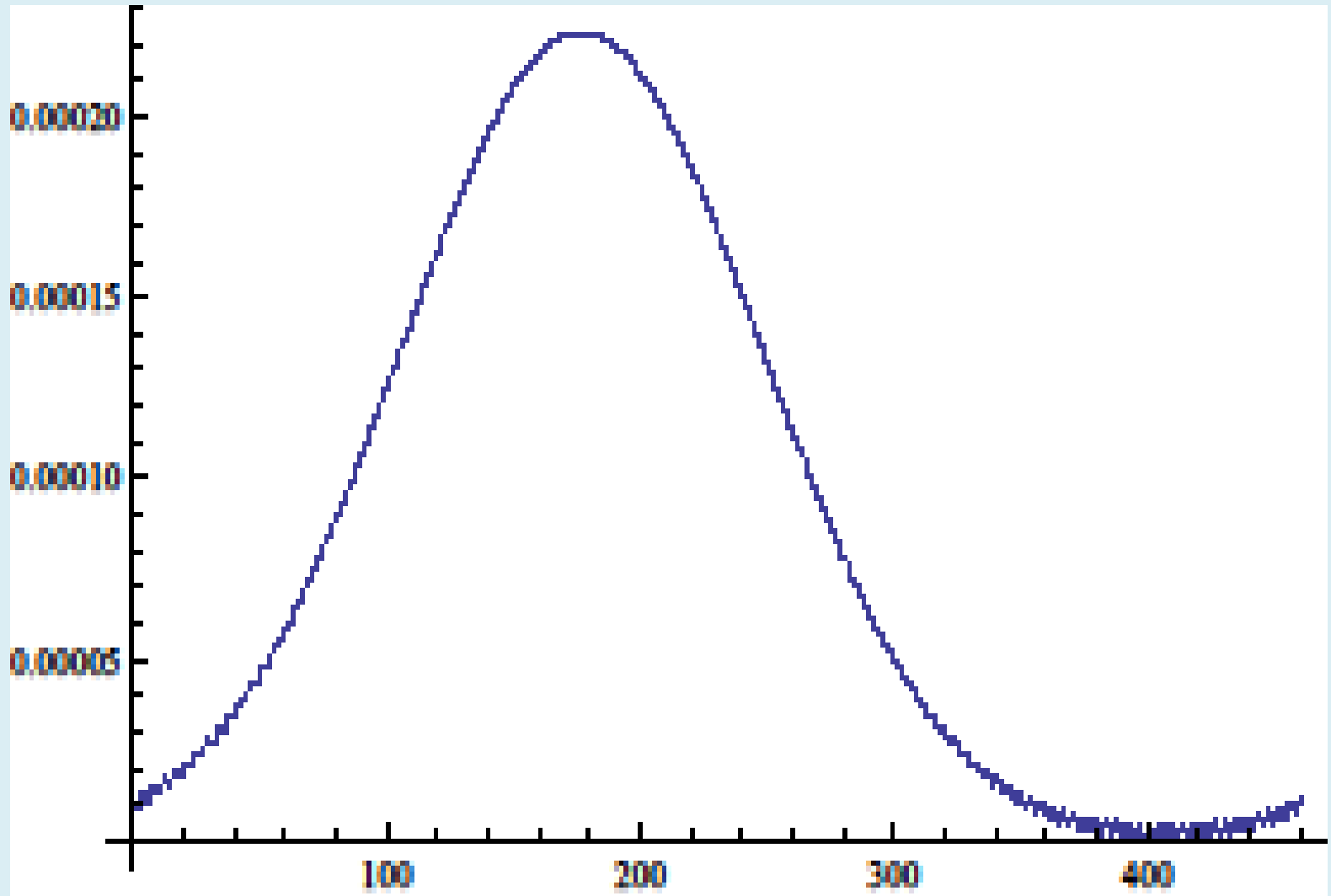
- TIME=1



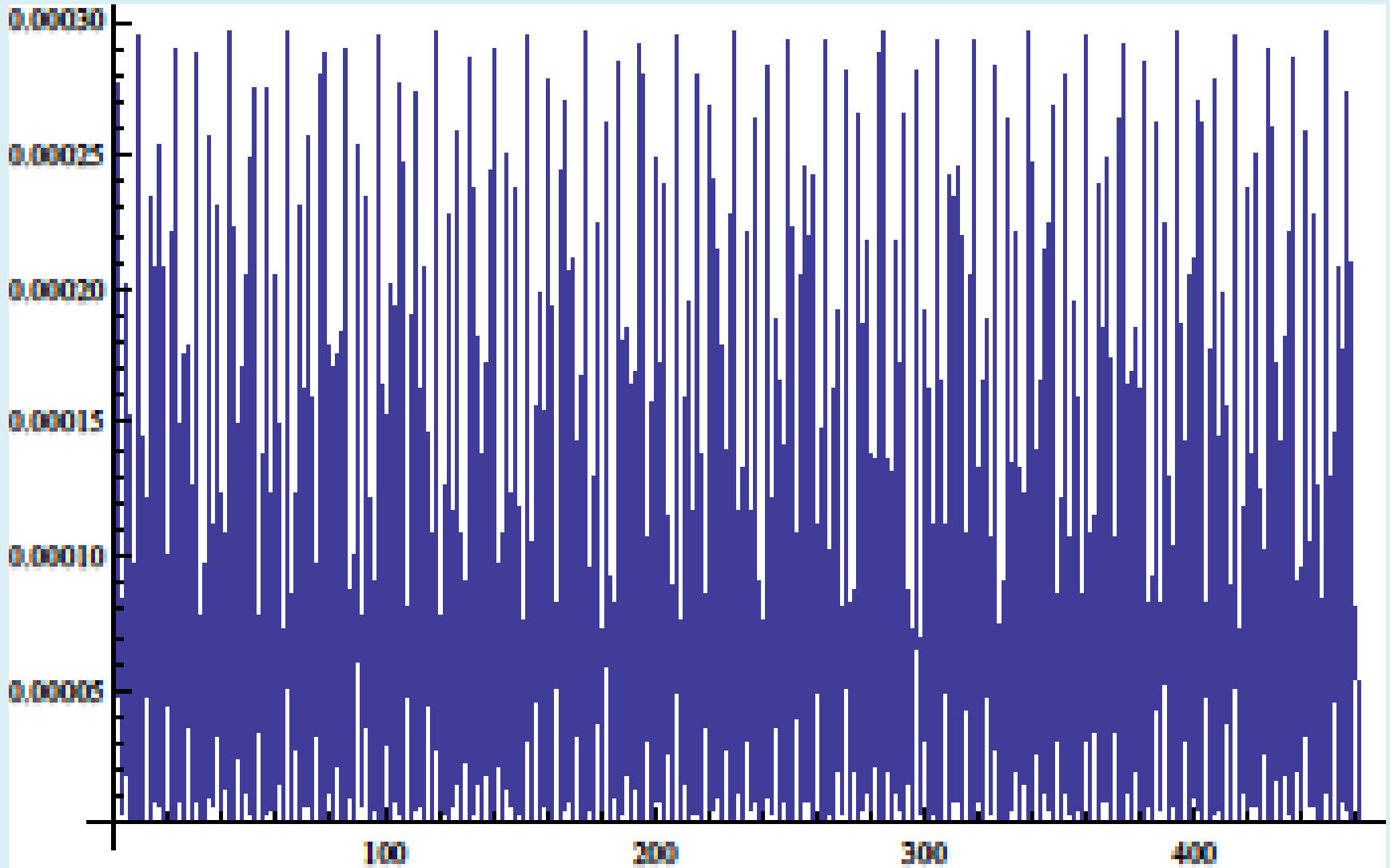
- $T=3$



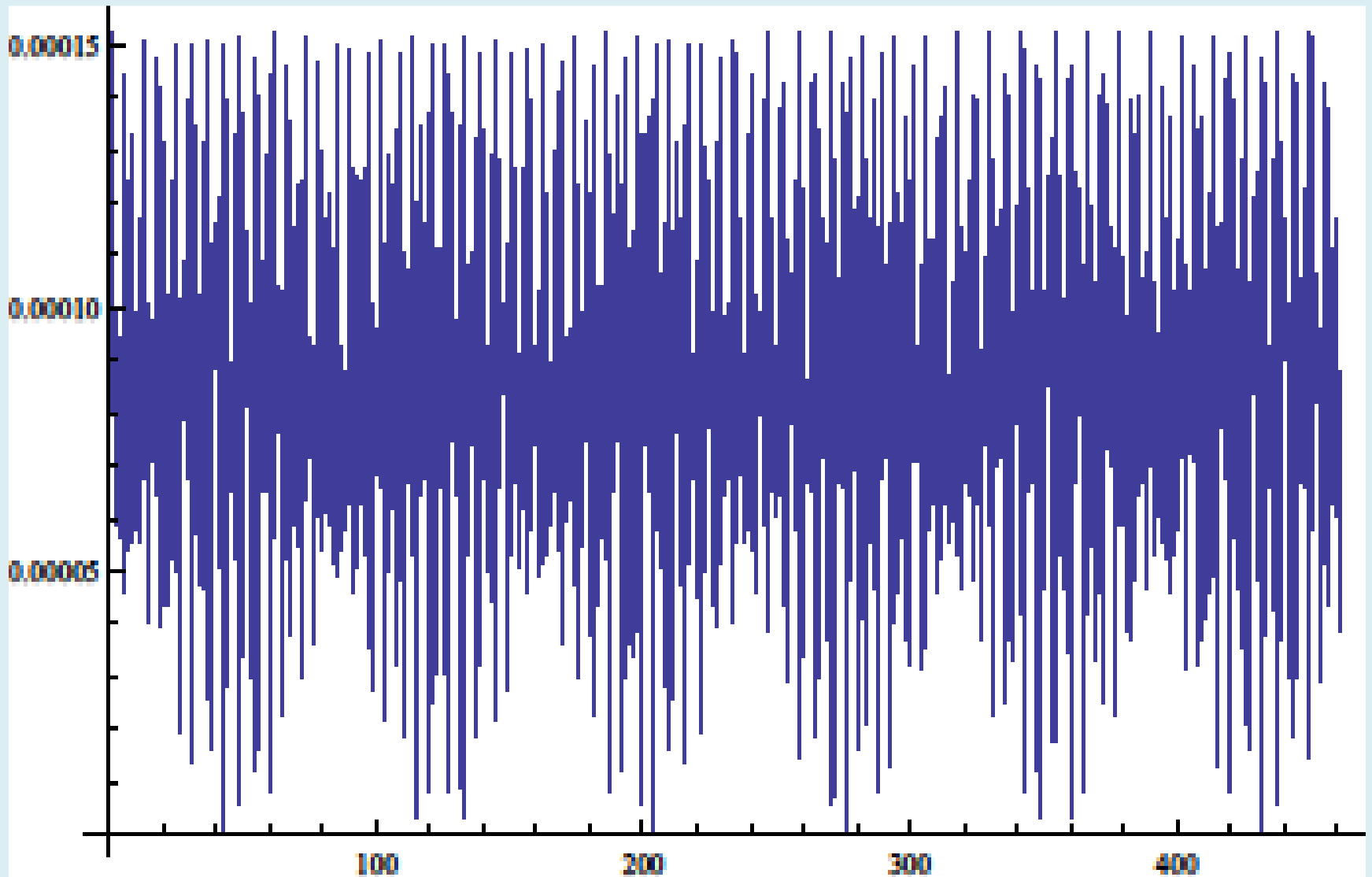
- $T=3$



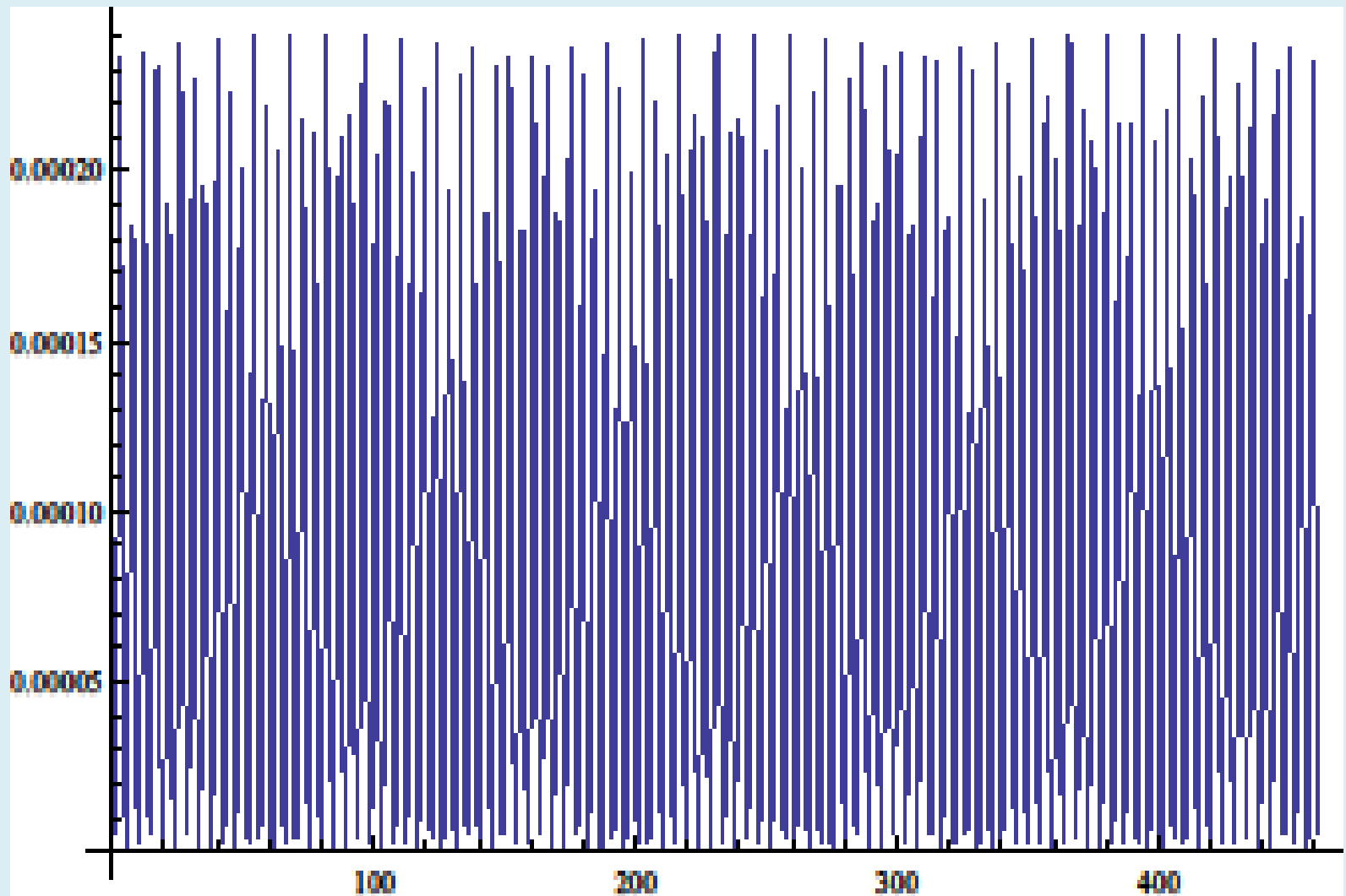
- $T=4$



- $T=6$



- $T=7$



CONCLUSIONS

DISCRETIZING AdS_2 RADIAL AND TIME NEAR HORIZON
GEOMETRY OF EXTREMAL BLACKHOLES
AND AT THE SAME TIME PRESERVING THE ALGEBRAIC
STRUCTURE OF THE ISOMETRIES
NECESSARILY LEADS TO THE $MOD[N]$ ARITHMETIC
DISCRETIZATION WITH HOLOGRAPHY $AdS_2[N]/CFT_1[N]$,
THIS CLASSICAL STRUCTURE IS LIFTED TO THE QUANTUM
LEVEL
THROUGH FINITE QUANTUM MECHANICS

THE QUANTUM CAT MAP (QACM) CHAOTIC DYNAMICS ON THE DISCRETIZED HORIZON $AdS_2[N]$,

DUE TO ITS RANDOM EIGENSTATES,(ETH),

THERMALIZES IN LOGARITHMIC TIME

SINGLE PARTICLE WAVE PACKETS,

THUS IT SATURATES

THE SCRAMBLING TIME BOUND OF HAYDEN-PRESKILL

SEKINO-SUSSKIND

THE COMPLEXITY OF THE QACM GROWS AS

$$n^2 = \text{Log}[S]^2$$