

Integrable deformations and non-commutativity

David Osten



Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

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based on 1608.08504 with Stijn van Tongeren
and ongoing work with Dieter Lüüst

Motivation

- integrability of toy models, for example in AdS/CFT:
type IIB GS superstring in $\text{AdS}_5 \times S^5 \leftrightarrow \mathcal{N} = 4$ SYM

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- integrability of toy models, for example in AdS/CFT:
type IIB GS superstring in $AdS_5 \times S^5 \leftrightarrow \mathcal{N} = 4$ SYM
- 'deformations' of integrable (2d) σ -models
 - AdS/CFT: less symmetric examples?
 - symmetries behind integrable structures?
 - generating new supergravity solutions?
 - non-abelian T -dualities and non-geometric backgrounds

Overview

- 1 Yang-Baxter deformations
Definition and properties
Overview
- 2 Connection to non-commutative physics
CFT-duals to Yang-Baxter deformed AdS-backgrounds
Embedding into a 'doubled space'
- 3 Conclusion

Yang-Baxter deformations - Definition and properties

- simplest setup: integrable deformations of principal chiral model on Lie group G ([Klimcik, 2008] & [Delduc et al., 2013])

$$S \propto \int d^2\sigma \operatorname{Tr} \left((g^{-1}\partial_+g) \frac{1}{\mathbb{1} - \eta R} (g^{-1}\partial_-g) \right)$$

for fields: $g : \Sigma \rightarrow G$ and with $\eta \in [0, 1)$.

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- Lax integrability (classical)
 - equations of motion equivalent to $\partial_+L_- - \partial_-L_+ - [L_+, L_-] = 0$, with $L_{\pm} \equiv L_{\pm}(\lambda) : \mathbb{C} \rightarrow \mathfrak{g}$
 \Rightarrow infinite tower of conserved charges
 - Deformed model is Lax integrable, if $R \in \operatorname{End}(\mathfrak{g})$ solution to (modified) classical Yang-Baxter equation:
 $[R(M), R(N)] - R([R(M), N] + [M, R(N)]) = C[M, N]$.

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- generalisation to symmetric space (super)coset models (e.g. superstring on $AdS_5 \times S^5$ possible)
- question: Is this deformation a supergravity solution?

Yang-Baxter deformations - Overview

- **abelian Yang-Baxter deformations**

[Osten and van Tongeren, 2016]

- based on R -operators with non-vanishing R^{ij} with $[t_i, t_j] = 0$
- equivalent to θ -shifts from the $O(d, d)$ T -duality group:
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- generalisation: **unimodular Yang-Baxter deformations:**

[Borsato and Wulff, 2016]

- based on unimodular R -operators: $R^{AB}[t_A, t_B] = 0$
- R unimodular \Leftrightarrow background is a supergravity solution.
- conjectured to be related to non-abelian T -duality transformations in general (shown for unimodular R -operators of AdS_5 resp. $SO(2, 4)$ in [Hoare and Tseytlin, 2016])

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- **η -deformations** [Delduc et al., 2013, Arutyunov et al., 2013]

- generated by Drinfel'd-Jimbo R -operators solving mCYBE
- quantum group deformations of isometry algebra
- not a supergravity solution
- *but*: Poisson-Lie T -dual to a supergravity solution (λ -deformation)

CFT-duals to Yang-Baxter deformed AdS-backgrounds

- Yang-Baxter deformation
 - closed string vs. open string picture [Seiberg and Witten, 1999]
 - \leftrightarrow Drinfel'd twist of underlying Hopf algebra structure [van Tongeren, 2015]
- abelian twist of S^5 of $AdS_5 \times S^5$:
 - e.g. deformation based on $(U(1))^3$ of $SO(6) \leftrightarrow$ deformation of \mathcal{R} -symmetry in $\mathcal{N} = 4$ SYM
 - represented by \star -product between the corr. scalars in Lagrangian, e.g.

$$\phi \star \psi = \exp(i\pi\gamma(q_1(\phi)q_2(\psi) - q_2(\phi)q_1(\psi))) \phi\psi,$$

where q_1, q_2 : charges of $U(1) \times U(1)$

- deformation of AdS_5 in $AdS_5 \times S^5 \leftrightarrow \mathcal{N} = 4$ SYM on non-commutative spacetime (Groenewold-Moyal \star -product with $\Theta^{AB} \propto R^{AB}$)

Embedding into a 'doubled space'

- simple setup of a double field theory on Lie groups
[Hull and Reid-Edwards, 2009, Hassler, 2017]
 - Drinfel'd double \mathcal{D} : $2d$ -dim. Lie group with compatible $O(d, d)$ -metric η
 - polarisations $\hat{=}$ bialgebra decompositions, $\mathfrak{d} = \mathfrak{g}_i \oplus \mathfrak{g}_i^*$, where $(\mathfrak{g}_i, \mathfrak{g}_i^*)$: dual pair of *maximally isotropic subalgebras* w.r.t. to η
 - A polarisation has to close under group multiplication!
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- Yang-Baxter deformations, generated by solutions R of cYBe
 - $\left[\begin{array}{l} \text{Principal chiral model} \\ \text{polarisation: } (\mathfrak{g}, (\mathfrak{u}(1))^d) \end{array} \right] \mapsto \left[\begin{array}{l} \text{Yang-Baxter def. model} \\ \text{polarisation: } (\mathfrak{g}, \mathfrak{g}^*) \end{array} \right]$
with Lie algebra structure on \mathfrak{g}^* : $\bar{f}_c^{ab} = f_{cd}^{[a} R^{b]d}$
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 - 'trivial β -shift', e.g. $[\beta, \beta]_S = 0$
- conjectural note on η -deformations:
 - non-trivial β -shift, $[\beta, \beta]_S \propto \mathcal{C}_{\mathfrak{g}}^3$
 - \rightarrow Q -/ R -flux background?
 - a unified description of η -/ λ -deformations in this way?
 - gerbe formulation needed?

Conclusion

Summary:

- overview over integrable deformations
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Open questions:

- non-abelian twists in AdS/CFT
- more elaborate double field theory on group manifolds
 - gerbe formulation
 - non-trivial fluxes
 - gauge algebra
- η -/ λ -deformation as non-geometric flux backgrounds

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Thank you for your attention!

BACKUP 1: Classical Yang-Baxter Equation (CYBE)

- Common form:

$$[r_{12}, r_{13}] + [r_{12}, r_{23}] + [r_{13}, r_{23}] = 0 \quad \text{for } r \in \mathfrak{g} \otimes \mathfrak{g}$$

Transition from a skew-symmetric r -matrix to a R -operator:

$$r = a \wedge b := \frac{1}{2}(a \otimes b - b \otimes a) \quad \rightarrow \quad R(M) := \text{Tr}_2(r \cdot (1 \otimes M))$$

- There are deformations based on solutions R of
 - CYBE: $[R(M), R(N)] - R([R(M), N] + [M, R(N)]) = 0$
 - mCYBE: $[R(M), R(N)] - R([R(M), N] + [M, R(N)]) = \pm[M, N]$

BACKUP 2: λ -deformation

Let G be a Lie group, and $g : \Sigma \rightarrow G$ with generators $\{t_a\}$ and structure constants $(f^c)_{ab}$. The action [Sfetsos, 2014]

$$S(g) = S_{WZW,k}(g) + \frac{\lambda}{\pi} \int d^2\sigma (g^{-1}\partial_+g)^a \left(\frac{1}{\mathbb{1} - \lambda \text{Ad}_g^{-1}} \right)_{ab} (g^{-1}\partial_+g)^b$$

describes a deformation of the WZW-model. The deformation parameter λ can be written in terms of the WZW-level k and the coupling of a principal chiral model κ is $\lambda = \frac{k^2}{\kappa^2 + k}$. The limits are

- $\lambda \rightarrow 0$: undeformed WZW model
- $k \rightarrow \infty$: non-abelian T -dual of principal chiral model on G

$$S_{NATD} = \frac{1}{\pi} \int d\sigma^2 \partial_+ \chi_a (\mathbb{1} - \chi_c(f^c))^{-1,ab} \partial_- \chi_b$$

- $k \ll \kappa^2$: perturbed WZW

$$S(g) = S_{WZW,k}(g) + \frac{k^2}{\pi\kappa^2} \int d\sigma^2 (g^{-1}\partial_+g)^a (g^{-1}\partial_+g)^a$$



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