

Signatures of extra dimensions in gravitational waves

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Direct detection of gravitational waves by LIGO and Virgo
Scientific Collab. [\[arXiv:1602.03837\]](#), [\[arXiv:1606.04855\]](#), [\[arXiv:1706.01812\]](#)

⇒ new observational tool to probe nature and test theories.

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↔ models beyond four-dimensional (4d) General Relativity

Here: test idea of having N **extra dimensions**: $D = 4 + N$.

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If \exists extra dimensions

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Many models with extra dimensions, from pheno. to
qu. grav.: large extra dimensions (ADD models),
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supergravities, string theories, M-theory...

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Variety of models: number, size, shape of extra dimensions...

Previous literature: typically very model dependent

⇒ here, remain as **generic** as possible.

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Most work on gravitational waves is about source: compute waveform for some emission. In 4d, governed by

$$\square_4 h_{\mu\nu} \approx T_{\mu\nu}^{(1)} + \text{gauge fixing.}$$

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Here: **away from source** (avoids model dependence).

Assume waves emitted (initial conditions), study **propagation**

\leftrightarrow corrections to $\square_4 h_{\mu\nu} = 0 + \textit{gauge fixing}$ due to extra dim.?

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D -dimensional General Relativity with cosmo. constant

\rightarrow derive gravitational wave equation and gauge fixing on generic background

\rightarrow split dimensions: $D = 4 + N \Rightarrow$ split equations

\leftrightarrow **modifications** of those on $h_{\mu\nu}$? Yes!

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In general, too complicated to read-off effect on wave

\leftrightarrow restrict background to Minkowski $\times \mathcal{M}_N$.

Minkowski: \checkmark for physical purposes; \mathcal{M}_N compact Ricci-flat.

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⇒ **Two effects**:

1. Breathing mode: new polarization mode in massless wave.
2. Additional (massive) waves of high frequencies.

↔ Observable in a near future?

General equations for gravitational waves

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In D dimensions

$$\text{General Relativity: } S = \frac{1}{2\kappa_D} \int d^D x \sqrt{|g_D|} (\mathcal{R}_D - 2\Lambda_D)$$
$$\hookrightarrow \text{Einstein equation: } \mathcal{R}_{MN} - \frac{2\Lambda_D}{D-2} g_{DMN} = 0$$

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Background + fluctuation: $g_{DMN} = g_{MN} + h_{MN}$

\hookrightarrow develop equation at 0th and 1st order:

$\mathcal{R}_{MN}^{(0)} - \frac{2\Lambda_D}{D-2} g_{MN} = 0$, $\mathcal{R}_{MN}^{(1)} - \frac{2\Lambda_D}{D-2} h_{MN} = 0$

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$$1^{\text{st}} \text{ order: } -\frac{1}{2} \square_D^{(0)} h_{MN} + \mathcal{R}^{(0)S}{}_{MNP} g^{PQ} h_{QS} + \nabla_{(M}^{(0)} \mathcal{G}_{N)} = 0$$

where $\mathcal{G}_N = \nabla_P^{(0)} g^{PQ} h_{QN} - \frac{1}{2} \nabla_N^{(0)} h_D$, with $h_D = g^{MN} h_{MN}$.

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de Donder (Lorenz) gauge fixing: $\mathcal{G}_N = 0$

$$-\frac{1}{2} \square_D^{(0)} h_{MN} + \mathcal{R}^{(0)S}{}_{MNP} g^{PQ} h_{QS} = 0$$

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Split into $4 + N$ dimensions

Background: $ds^2 = e^{2A(y)} \tilde{g}_{\mu\nu}(x) dx^\mu dx^\nu + g_{mn}(y) dy^m dy^n$

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h_{MN} : $h_{\mu\nu}, h_{\mu m}, h_{mn}$, generic coordinate dependence

traces $\tilde{h}_4 = h_{\mu\nu} \tilde{g}^{\nu\mu}$, $h_4 = h_{\mu\nu} \tilde{g}^{\nu\mu} e^{-2A}$, $h_N = h_{mn} g^{nm}$.

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D-dimensional wave equation: MN components:

$$\begin{aligned} & e^{-2A}\tilde{\alpha}_4\mathbf{h}_{\mu\nu} + \Delta_{\mathcal{M}}h_{\mu\nu} - h_{\mu\nu}\Delta_{\mathcal{M}}\ln e^{2A} \\ & - 2\tilde{\mathcal{R}}^\pi{}_{\mu\nu\sigma}g^{\sigma\rho}h_{\rho\pi} - \frac{1}{2}e^{-2A}g^{pq}\partial_p e^{2A}\partial_q e^{2A}\left(\tilde{g}_{\nu\mu}h_4 - h_{\nu\mu}e^{-2A}\right) \\ & - e^{-2A}\tilde{\nabla}_{(\mu}h_{\nu)m}g^{mn}\partial_n e^{2A} - \tilde{g}_{\mu\nu}h^{rp}\left(\nabla_r\partial_p e^{2A} + \frac{1}{2}e^{-2A}\partial_r e^{2A}\partial_p e^{2A}\right) = 0 \\ & e^{-2A}\tilde{\alpha}_4\mathbf{h}_{\mu n} + \Delta_{\mathcal{M}}h_{\mu n} + e^{-2A}g^{pq}\nabla_p h_{\mu n}\partial_q e^{2A} + e^{-2A}h_{\mu m}g^{mp}\nabla_n\partial_p e^{2A} \\ & - 2e^{-4A}h_{\mu m}g^{mp}\partial_p e^{2A}\partial_n e^{2A} - e^{-4A}h_{\mu n}g^{pq}\partial_p e^{2A}\partial_q e^{2A} - \frac{1}{2}h_{\mu n}\Delta_{\mathcal{M}}\ln e^{2A} \\ & - e^{-4A}\tilde{g}^{\pi\rho}\tilde{\nabla}_\pi h_{\mu\rho}\partial_n e^{2A} + e^{-2A}g^{pq}\partial_\mu h_{np}\partial_q e^{2A} = 0 \\ & e^{-2A}\tilde{\alpha}_4\mathbf{h}_{mn} + \Delta_{\mathcal{M}}h_{mn} + 2e^{-2A}g^{pq}\partial_p e^{2A}\nabla_q h_{mn} + 2g^{pq}\partial_p e^{-2A}h_{q(m}\partial_n)e^{2A} \\ & - 2\mathcal{R}^s{}_{mnp}g^{pq}h_{qs} - 2e^{-4A}\tilde{g}^{\pi\rho}\tilde{\nabla}_\pi h_{\rho(m}\partial_n)e^{2A} - h_4\nabla_n(e^{-2A}\partial_m e^{2A}) = 0 \end{aligned}$$

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 & - 2\tilde{\mathcal{R}}^\pi{}_{\mu\nu\sigma\rho}g^{\sigma\rho}h_{\rho\pi} - \frac{1}{2}e^{-2A}g^{pq}\partial_p e^{2A}\partial_q e^{2A}\left(\tilde{g}_{\nu\mu}h_4 - h_{\nu\mu}e^{-2A}\right) \\
 & - e^{-2A}\tilde{\nabla}_{(\mu}h_{\nu)m}g^{mn}\partial_n e^{2A} - \tilde{g}_{\mu\nu}h^{rp}\left(\nabla_r\partial_p e^{2A} + \frac{1}{2}e^{-2A}\partial_r e^{2A}\partial_p e^{2A}\right) = 0 \\
 & e^{-2A}\tilde{\square}_4\mathbf{h}_{\mu n} + \Delta_{\mathcal{M}}h_{\mu n} + e^{-2A}g^{pq}\nabla_p h_{\mu n}\partial_q e^{2A} + e^{-2A}h_{\mu m}g^{mp}\nabla_n\partial_p e^{2A} \\
 & - 2e^{-4A}h_{\mu m}g^{mp}\partial_p e^{2A}\partial_n e^{2A} - e^{-4A}h_{\mu n}g^{pq}\partial_p e^{2A}\partial_q e^{2A} - \frac{1}{2}h_{\mu n}\Delta_{\mathcal{M}}\ln e^{2A} \\
 & - e^{-4A}\tilde{g}^{\pi\rho}\tilde{\nabla}_\pi h_{\mu\rho}\partial_n e^{2A} + e^{-2A}g^{pq}\partial_\mu h_{np}\partial_q e^{2A} = 0 \\
 & e^{-2A}\tilde{\square}_4\mathbf{h}_{mn} + \Delta_{\mathcal{M}}h_{mn} + 2e^{-2A}g^{pq}\partial_p e^{2A}\nabla_q h_{mn} + 2g^{pq}\partial_p e^{-2A}h_{q(m}\partial_n)e^{2A} \\
 & - 2\mathcal{R}^s{}_{mnp}g^{pq}h_{qs} - 2e^{-4A}\tilde{g}^{\pi\rho}\tilde{\nabla}_\pi h_{\rho(m}\partial_n)e^{2A} - h_4\nabla_n(e^{-2A}\partial_m e^{2A}) = 0
 \end{aligned}$$

D-dimensional de Donder gauge:

$$\begin{aligned}
 & e^{-2A}\tilde{g}^{\pi\rho}\tilde{\nabla}_\pi h_{\rho\nu} - \frac{e^{-2A}}{2}\tilde{\nabla}_\nu\tilde{h}_4 - \frac{1}{2}\nabla_\nu h_N + \nabla^q h_{q\nu} + 2h_{p\nu}g^{pq}e^{-2A}\partial_q e^{2A} = 0 \\
 & g^{pq}\nabla_p h_{qr} - \frac{e^{-2A}}{2}\nabla_r\tilde{h}_4 - \frac{1}{2}\nabla_r h_N + g^{\pi\rho}\tilde{\nabla}_\pi h_{\rho r} + 2h_{mr}g^{mp}e^{-2A}\partial_p e^{2A} = 0
 \end{aligned}$$

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$$\begin{aligned} & e^{-2A} \tilde{\square}_4 h_{\mu\nu} + \Delta_{\mathcal{M}} h_{\mu\nu} - h_{\mu\nu} \Delta_{\mathcal{M}} \ln e^{2A} \\ & - 2 \tilde{\mathcal{R}}^{\pi}_{\mu\nu\sigma} g^{\sigma\rho} h_{\rho\pi} - \frac{1}{2} e^{-2A} g^{pq} \partial_p e^{2A} \partial_q e^{2A} \left(\tilde{g}_{\nu\mu} h_4 - h_{\nu\mu} e^{-2A} \right) \\ & - e^{-2A} \tilde{\nabla}_{(\mu} h_{\nu)m} g^{mn} \partial_n e^{2A} - \tilde{g}_{\mu\nu} h^{rp} \left(\nabla_r \partial_p e^{2A} + \frac{1}{2} e^{-2A} \partial_r e^{2A} \partial_p e^{2A} \right) = 0 \end{aligned}$$

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Many terms, coupling to $h_{\mu n}$ and h_{mn} , or $\partial_p e^{2A}$.

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\hookrightarrow **constant** e^{2A} : $\partial_p e^{2A} = 0$. $e^{2A} = 1$, $\tilde{g}_{\mu\nu} \rightarrow g_{\mu\nu}$.

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For physics: **Minkowski**

background equation \Rightarrow Ricci-flat \mathcal{M}_N : $\mathcal{R}_{mn} = 0$ (e.g. any CY).

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 & e^{-2A} \tilde{\square}_4 h_{\mu\nu} + \Delta_{\mathcal{M}} h_{\mu\nu} - h_{\mu\nu} \Delta_{\mathcal{M}} \ln e^{2A} \\
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$$\square_4 h_{\mu n} + \Delta_{\mathcal{M}} h_{\mu n} = 0$$

$$\square_4 h_{mn} + \Delta_{\mathcal{M}} h_{mn} = 2\mathcal{R}^s_{mnp} g^{pq} h_{qs}$$

$$g^{\pi\rho} \nabla_{\pi} h_{\rho\nu} - \frac{1}{2} \nabla_{\nu} h_4 - \frac{1}{2} \nabla_{\nu} h_N + g^{pq} \nabla_p h_{q\nu} = 0$$

$$g^{\pi\rho} \nabla_{\pi} h_{\rho r} - \frac{1}{2} \nabla_r h_4 - \frac{1}{2} \nabla_r h_N + g^{pq} \nabla_p h_{qr} = 0$$

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$$g^{mn} \times \square_4 h_{mn} + \Delta_{\mathcal{M}} h_{mn} = 2\mathcal{R}^s_{mnp} g^{pq} h_{qs}$$

$$\hookrightarrow \square_4 h_N + \Delta_{\mathcal{M}} h_N = 0$$

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Equation analysis

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Interested in 4d wave $h_{\mu\nu}$, e.g. in $\square_4 h_{\mu\nu} + \Delta_{\mathcal{M}} h_{\mu\nu} = 0$

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Consider \mathcal{M}_N compact (without boundary)

→ use basis of eigenfunctions $\{\omega_{\mathbf{k}}(y)\}$ of $\Delta_{\mathcal{M}}$,

discrete basis, label \mathbf{k} : $\Delta_{\mathcal{M}} \omega_{\mathbf{k}} = -m_{\mathbf{k}}^2 \omega_{\mathbf{k}}$ (e.g. T^N : $\omega_{\mathbf{k}}(y) = e^{i\mathbf{k}\cdot\mathbf{y}}$)

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David
ANDRIOT

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“Coupling” with zero-mode of internal trace $h_N^0 = (g^{mn} h_{mn})^0$.

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Both decouple from other fields/modes → analyse this system.

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With these equations: residual gauge freedom

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Fourier expansion on plane waves with wave vector k^ρ :

$$h_{\mu\nu}^0 = \int d^4k e_{\mu\nu}^k \operatorname{Re}\{e^{ik_\rho x^\rho}\}, \quad h_N^0 = \int d^4k f_N^k \operatorname{Re}\{e^{ik_\rho x^\rho}\}$$

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On each plane-wave, gauge condition + residual gauge freedom

$$e_{ij}^k = \left(\begin{array}{ccc} e_{11}^k & e_{12}^k & 0 \\ e_{12}^k & -e_{11}^k - f_N^k & 0 \\ 0 & 0 & 0 \end{array} \right)_{ij}$$

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\hookrightarrow G. R. h_{ij}^\times, h_{ij}^+ polarization modes and **breathing mode** h_{ij}^0 .

Massive modes

Focus on $h_{\mu\nu}^{\mathbf{k}\neq\mathbf{0}}$: equations: $\square_4 h_{\mu\nu}^{\mathbf{k}} - m_{\mathbf{k}}^2 h_{\mu\nu}^{\mathbf{k}} = 0$ + gauge cond.

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Residual gauge freedom \Rightarrow fix it (subtle)

\Rightarrow standard Transverse-Traceless massive graviton:

$$\partial^\nu h_{\mu\nu}^{\mathbf{k}} = 0, \quad h_4^{\mathbf{k}} = 0$$

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All six polarization modes, only 5 independent ones.

Massive modes

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All six polarization modes, only 5 independent ones.

(High) angular frequency $\omega_{\mathbf{k}} \sim m_{\mathbf{k}}$.

Two effects and observability

1. New polarization mode in massless wave: **breathing mode**.
2. **Additional** (massive) **waves** of high frequencies.

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Breathing mode in the massless wave

Each polarization mode \rightarrow **specific space deformation**

(stretch and shrink) with $\xi^i = x_0^i + \Delta x^i$

Geodesic equation $\ddot{\xi}_i = \frac{1}{2} \ddot{h}_{ij}^0 \xi^j \rightsquigarrow \Delta x^i = \frac{1}{2} h_{ij}^0 x_0^j$

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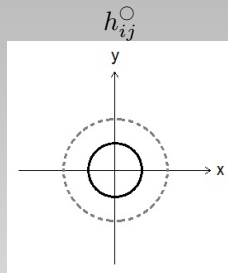
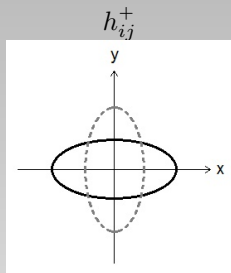
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Deformation of test-point circle in transverse plane:



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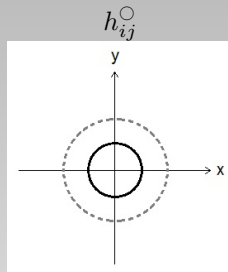
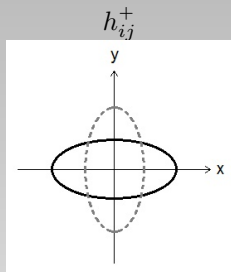
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Breathing mode: need several detectors, different orientations

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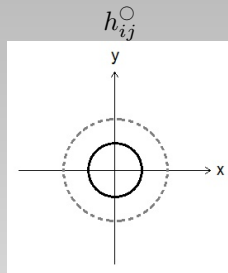
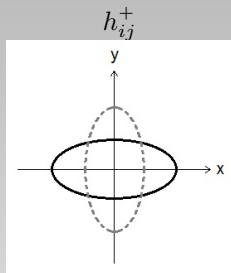
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Breathing mode: need several detectors, different orientations
Amplitude? Related to that h_N^0 ... Emission?

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Additional (massive) waves

All six polarization modes \rightarrow various space deformations.

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Worse in future (planned) detectors.

Energy \rightarrow amplitude is low...

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Thank you for your attention!